

Neutrinoless double-delta decay and The MAJORANA DEMONSTRATOR

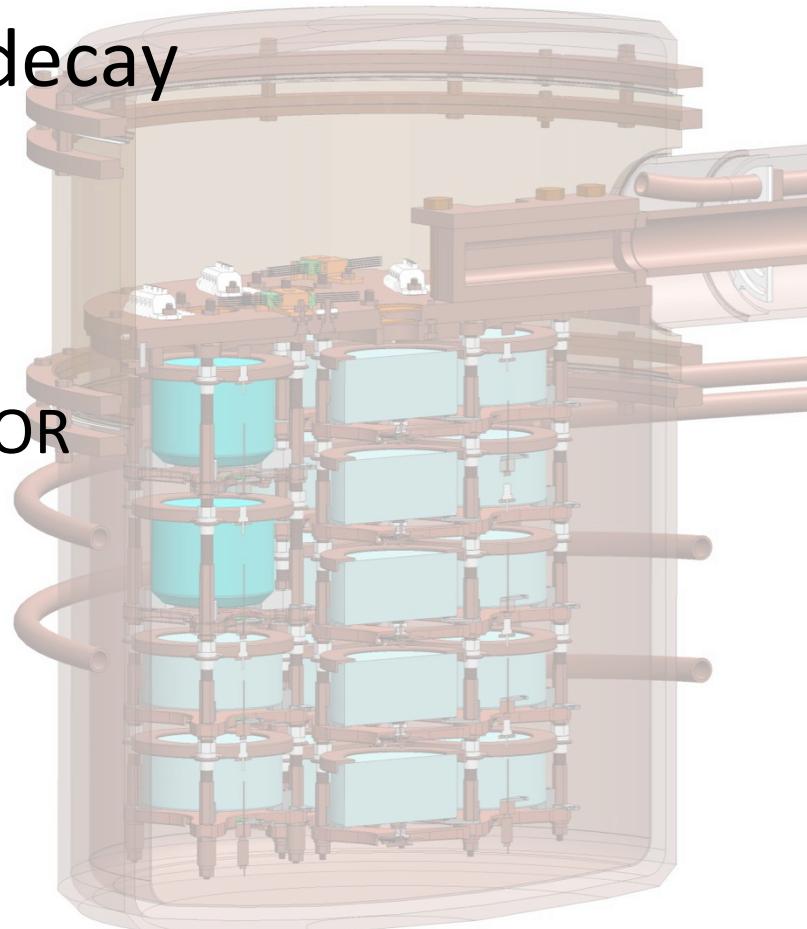
Wenqin Xu, LANL

For the MAJORANA collaboration

Santa Fe Summer Workshop
Implications of Neutrino Flavor Oscillations
INFO 2013

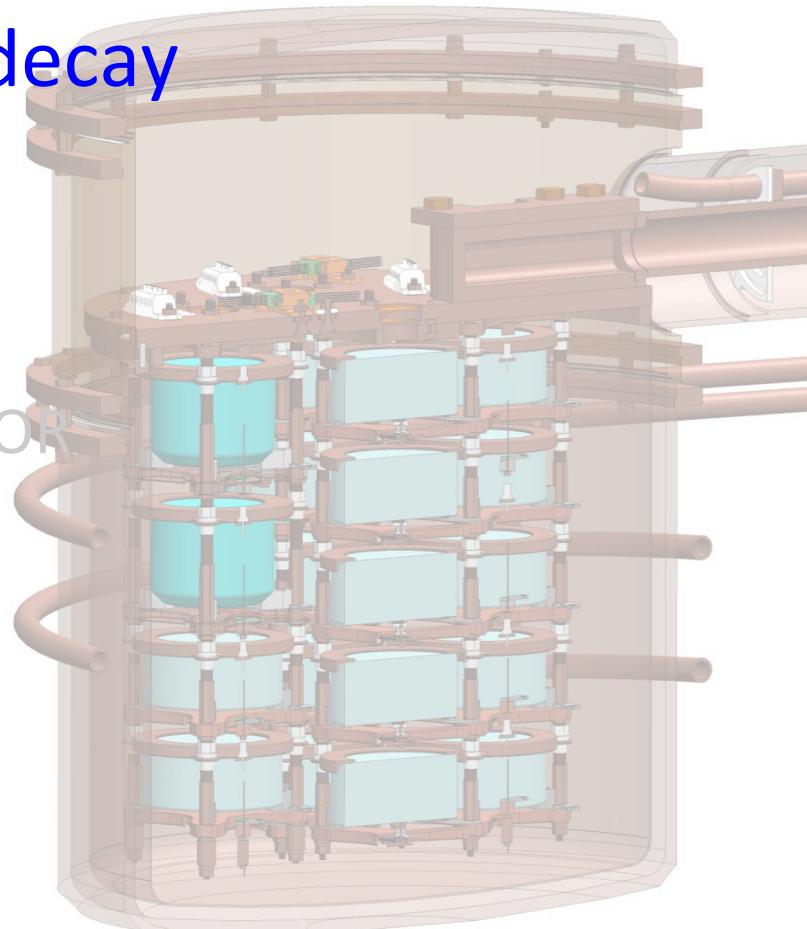
Outline of the talk

- Neutrinoless double-beta decay
 - the physics
 - the experiments
- The MAJORANA DEMONSTRATOR



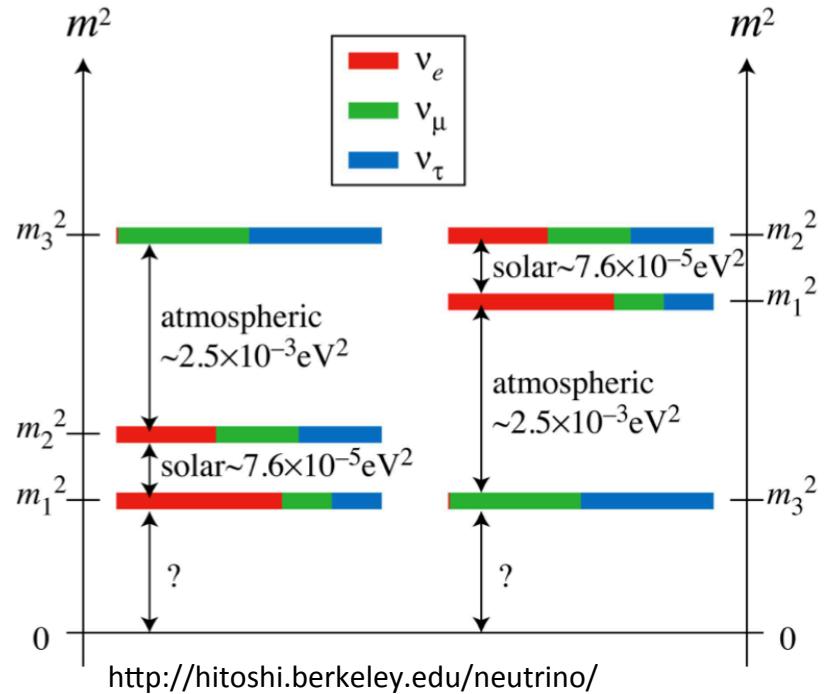
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Questions for neutrino physics

- Neutrinoless double beta decay:**
- Is neutrino its own antiparticle (i.e. Majorana Particle)?
 $0\nu\beta\beta$ is the only practical way to test this.
 - Is lepton number violated?
 - Leptogenesis as a way to produce the excess of matter?
 - Neutrino mass hierarchy ?
 - Absolute neutrino mas scale?

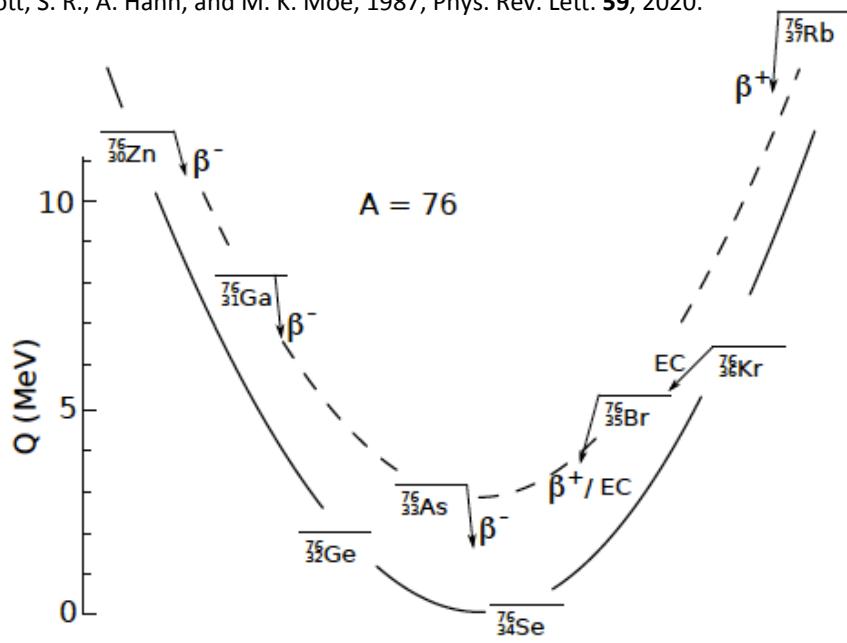


2-neutrino double beta decay

First direct observation by Steve Elliott et al in 1987

Elliott, S. R., A. Hahn, and M. K. Moe, 1987, Phys. Rev. Lett. **59**, 2020.

2nd order weak decay, very long half life

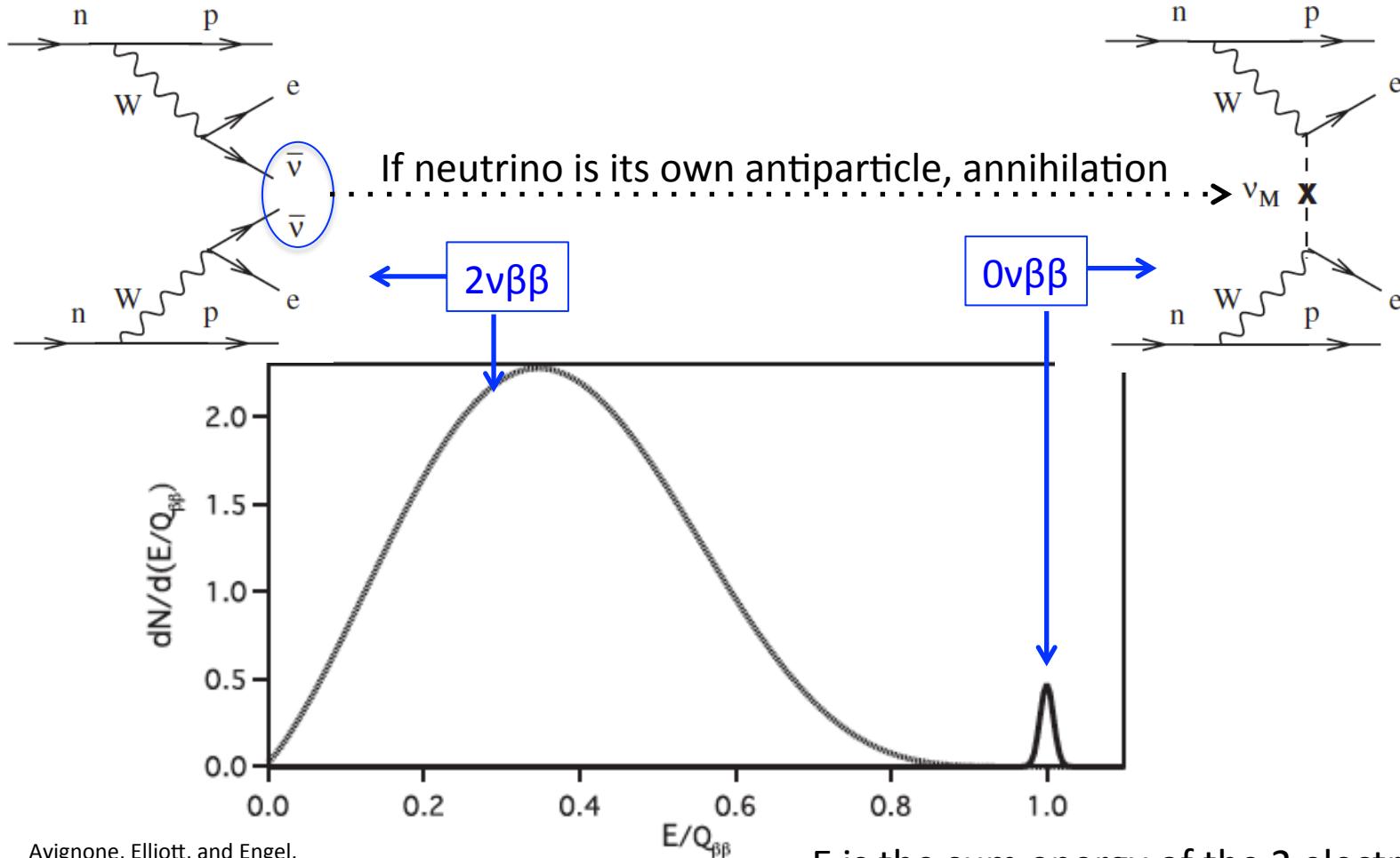


- Only possible if single beta decay is energetically forbidden.
- Observed for some nuclei with even numbers of protons and neutrons

Isotope	$T_{1/2}(2\nu)$, yr
^{48}Ca	$4.3^{+2.1}_{-1.0} \times 10^{19}$
^{76}Ge	$(1.5 \pm 0.1) \times 10^{21}$
^{82}Se	$(0.92 \pm 0.07) \times 10^{20}$
^{96}Zr	$(2.0 \pm 0.3) \times 10^{19}$
^{100}Mo	$(7.1 \pm 0.4) \times 10^{18}$
$^{100}\text{Mo}-^{100}\text{Ru}(0_1^+)$	$(6.2^{+0.9}_{-0.7}) \times 10^{20}$
^{116}Cd	$(3.0 \pm 0.2) \times 10^{19}$
^{128}Te	$(2.5 \pm 0.3) \times 10^{24}$
^{130}Te	$(0.9 \pm 0.1) \times 10^{21}$
^{150}Nd	$(7.8 \pm 0.7) \times 10^{18}$
$^{150}\text{Nd}-^{150}\text{Sm}(0_1^+)$	$1.4^{+0.5}_{-0.4} \times 10^{20}$
^{238}U	$(2.0 \pm 0.6) \times 10^{21}$
$^{130}\text{Ba}; \text{ECEC}(2\nu)$	$(2.2 \pm 0.5) \times 10^{21}$

A. S. Barabash, ISSN 1063-7788, Physics of Atomic Nuclei, 2010, Vol. 73, No. 1, pp. 162–178.

Neutrinoless double beta decay ($0\nu\beta\beta$)



Avignone, Elliott, and Engel,
Rev. Mod. Phys., Vol. 80, No. 2, April–June 2008

$0\nu\beta\beta$ half life is related to mass

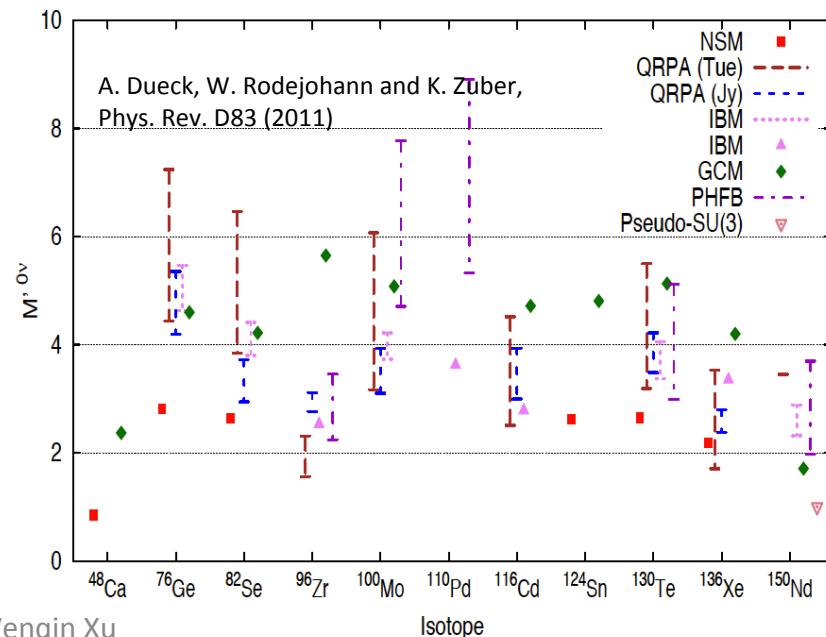
$$(T_{1/2}^{0\nu})^{-1} = G^{0\nu} |M_{0\nu}|^2 \left(\frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

$\langle m_{\beta\beta} \rangle = \left| \sum_{i=1} U_{ei}^2 m_i \right|$: Effective Majorana neutrino mass

$G^{0\nu}$: Phase factor

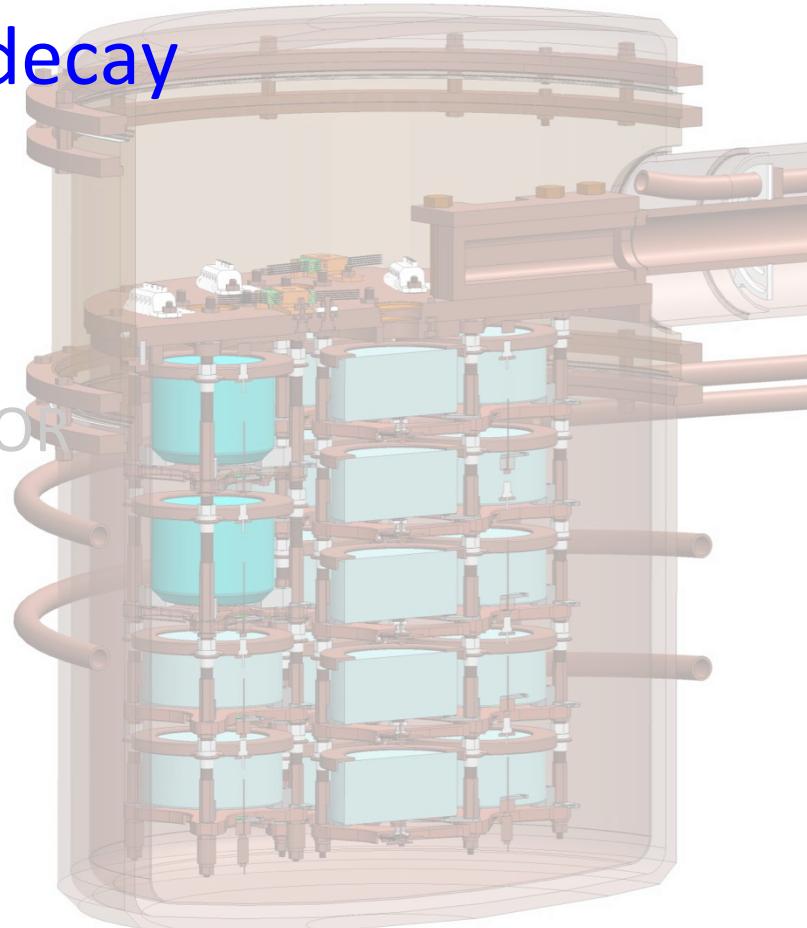
$M_{0\nu}$: Nuclear Matrix Element

Half life can be directly translated to effective majorana neutrino mass, although large uncertainties exist on NME calculation



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Recent List of $0\nu\beta\beta$ experiments

Isotope	$G^{0\nu}$ $\left[\frac{10^{-14}}{\text{yr}} \right]$	$Q_{\beta\beta}$ [keV]	Nat. ab. [%]	$T_{1/2}^{2\nu}$ $[10^{20} \text{ yr}]$	Experiments
^{48}Ca	6.3	4273.7	0.187	0.44	CANDLES
^{76}Ge	0.63	2039.1	7.8	15	GERDA, MAJORANA DEMONSTR.
^{82}Se	2.7	2995.5	9.2	0.92	SuperNEMO, Lucifer
^{100}Mo	4.4	3035.0	9.6	0.07	MOON, AMoRe
^{116}Cd	4.6	2809.1	7.6	0.29	Cobra
^{130}Te	4.1	2530.3	34.5	9.1	CUORE
^{136}Xe	4.3	2457.8	8.9	21	EXO, Next, Kamland-Zen
^{150}Nd	19.2	3367.3	5.6	0.08	SNO+, DCBA/MTD

B. Schwingenheuer, Ann. Phys. (Berlin) 525, No. 4 (2013)

The Choice of Ge

[Steven R. Elliott, Petr Vogel, Ann.Rev.Nucl.Part.Sci.52:115-151,2002](#)

Excellent energy resolution: crucial in distinguishing
2νββ (the ultimate background) from 0νββ near the end point

$F = \frac{7Q\delta^6}{m_e}$ is roughly the fraction of 2νββ decays ends up in the 0νββ peak region, $\delta = \Delta E/Q$

$\frac{S}{B} = \frac{m_e}{7Q\delta^6} \frac{\Gamma_{0\nu}}{\Gamma_{2\nu}} = \frac{m_e}{7Q\delta^6} \frac{T_{1/2}^{2\nu}}{T_{1/2}^{0\nu}}$. depends on the half life of 2νββ and 0νββ, element dependent

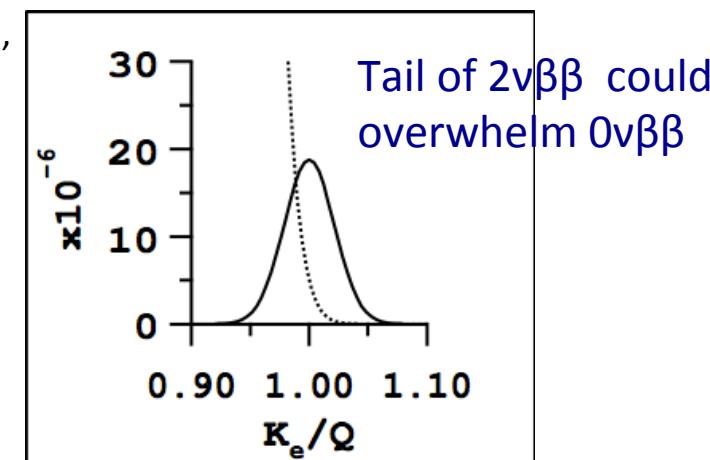
Half life ratio and $(\Delta E)^6$ decides the S/B, $\Gamma_{0\nu}=1e-6 \Gamma_{2\nu}$ →
thus the ultimate sensitivity.

ΔE is crucial

Ge detector has $\Delta E/E \sim 0.2\%$ at $Q_{\beta\beta}$

i.e. 4 keV Region of Interest (ROI) @ 2039 keV

2νββ is estimated to be negligible of the final
background in ROI for the DEMONSTRATOR



Good reasons for Ge

- ✓ Excellent Energy resolution ($\sim 0.2\%$ at 2039keV)
- ✓ Source is detector
- ✓ Can be enriched in ^{76}Ge to 86%
- ✓ Low level of radio-impurities can be achieved during processing
- ✓ Technology is well understood
- ✓ Easy to operate (LN temperature, volume is small)
- ✓ Large Q-value puts $0\nu\beta\beta$ peak above most backgrounds

Previous Ge experiment

The upper limits on $0\nu\beta\beta$ half-life:

Heidelberg-Moscow:

$T_{1/2}({}^{76}\text{Ge}) > 1.9 \times 10^{25}$ years (90% CL)

Eur. Phys. J. A. 12, 147-154 (2001)

IGEX (International Germanium Experiment):

$T_{1/2}({}^{76}\text{Ge}) > 1.57 \times 10^{25}$ years (90% CL)

Phys. Rev. D, 65, 092007 (2002)

They are the most sensitive limits until recently.

Recent Non-Ge experiments:

EXO-200: $T_{1/2}({}^{136}\text{Xe}) > 1.6 \times 10^{25}$ years (90% CL)

Phys. Rev. Lett. 109, 032505 (2012)

and

KamLAND-Zen: $T_{1/2}({}^{136}\text{Xe}) > 1.9 \times 10^{25}$ years (90% CL)

Phys. Rev. Lett. 110, 062502 (2013)

The (2004) 4.2σ $0\nu\beta\beta$ claim: $T_{1/2} = 1.19^{+0.38}_{-0.22} \times 10^{25}$ yr

Total exposure $71.8 \text{ kg}^*\text{yr}$

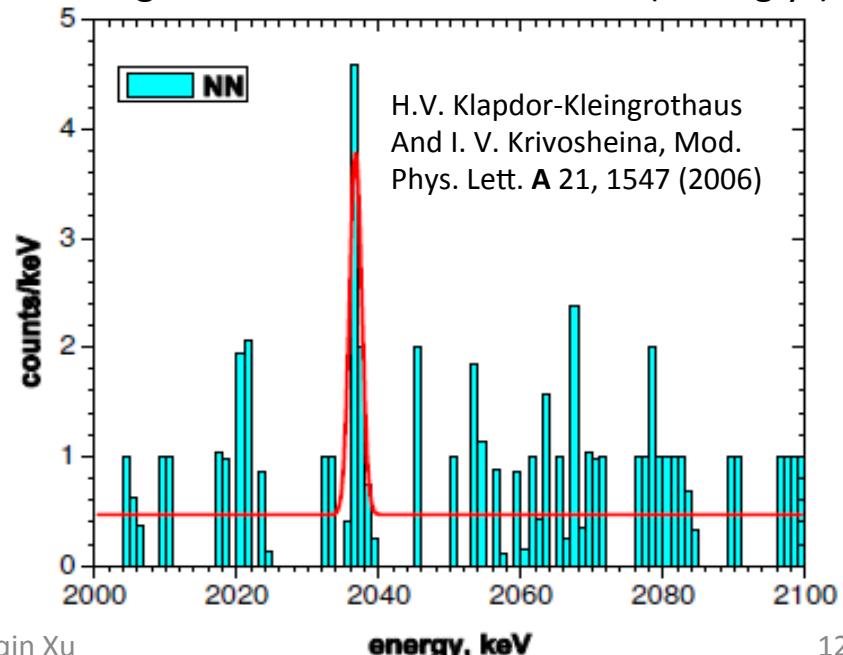
H.V. Klapdor-Kleingrothaus, et al, PLB 586 (2004) 198-212

The (2006) 6.4σ $0\nu\beta\beta$ claim: $T_{1/2} = 2.23^{+0.44}_{-0.31} \times 10^{25}$ yr

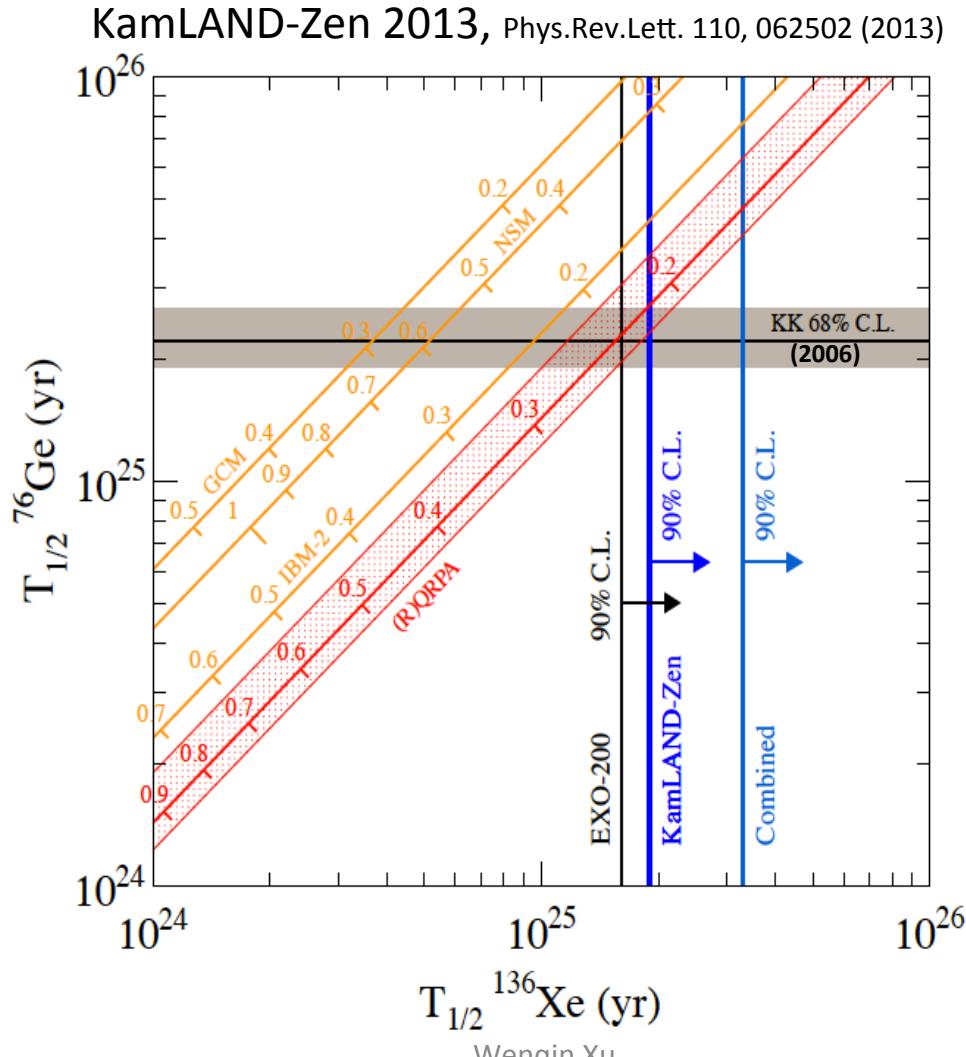
Reanalysis with Pulse Shape Decimation (PSA) used

Double-peak structure reported at $Q_{\beta\beta}$

Low background after PSD: $0.015 \text{ cnts}/(\text{keV}\cdot\text{kg}\cdot\text{yr})$



Status prior to GERDA and MJD

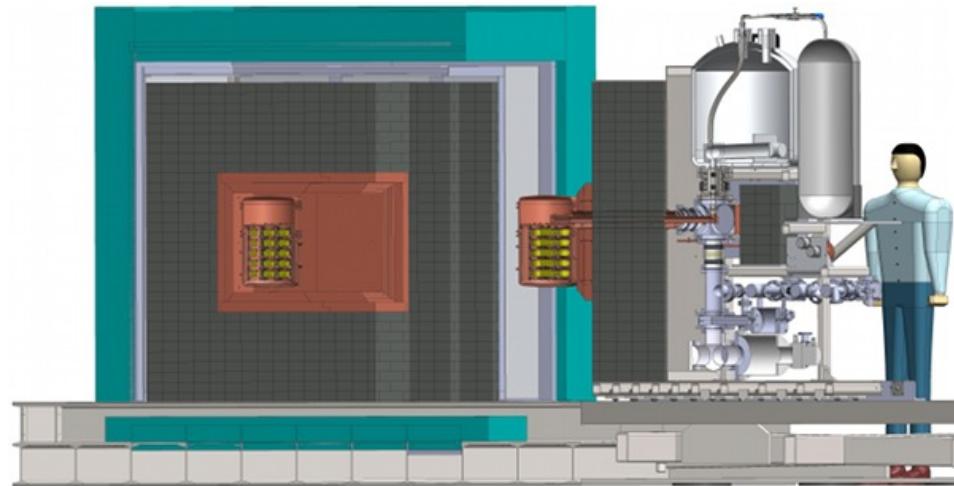


GERDA & MAJORANA DEMONSTRATOR



GERDA

Detectors inside a liquid argon shield
high-Z material budget is small,
relaxing depth requirement



MAJORANA DEMONSTRATOR

Detectors are inside layers of compact shield
Effectively shield against natural radioactivity

GERDA: experimental setup

Eur. Phys. J. C (2013) 73:2330
[arXiv:1212.4067](https://arxiv.org/abs/1212.4067)

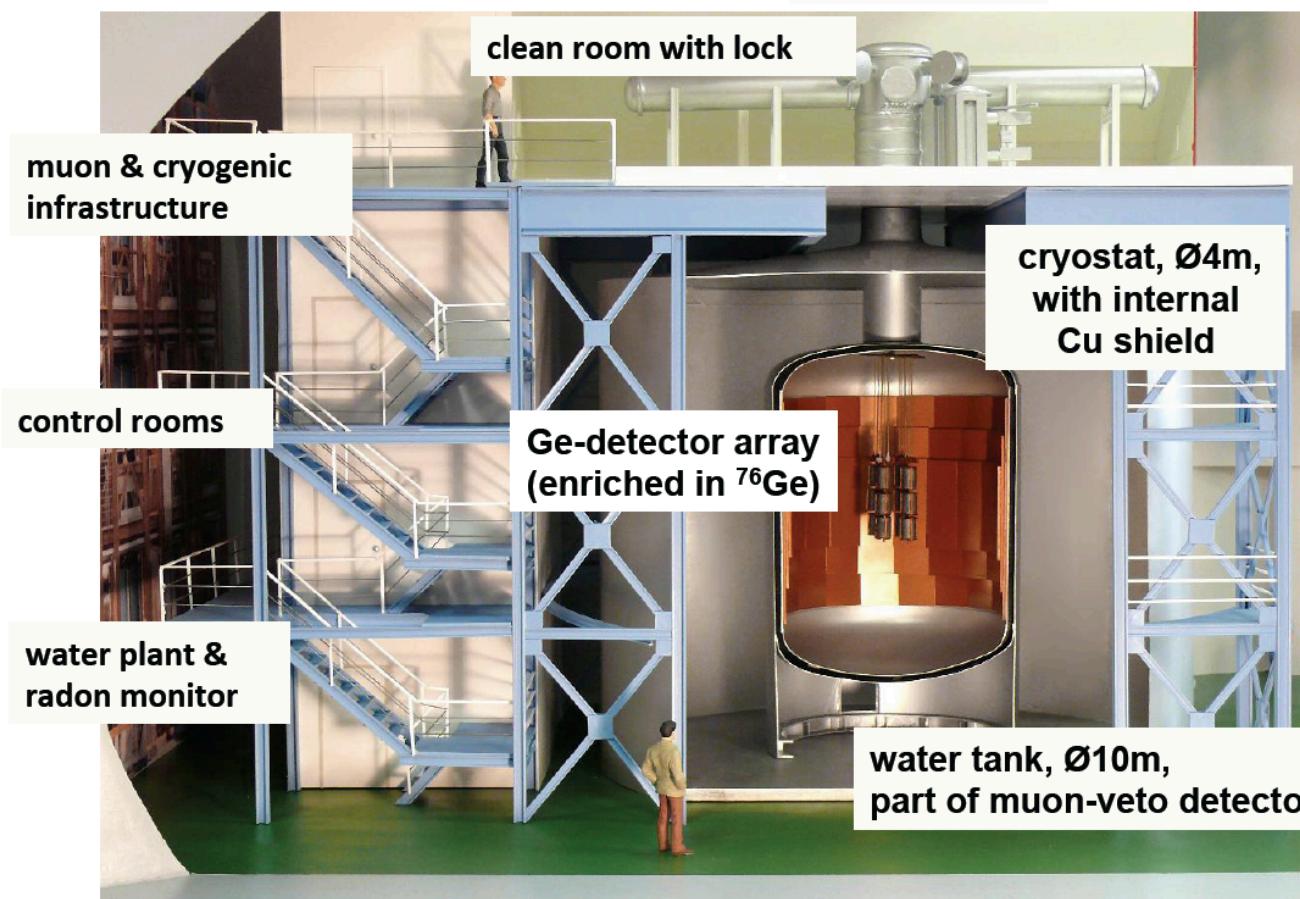
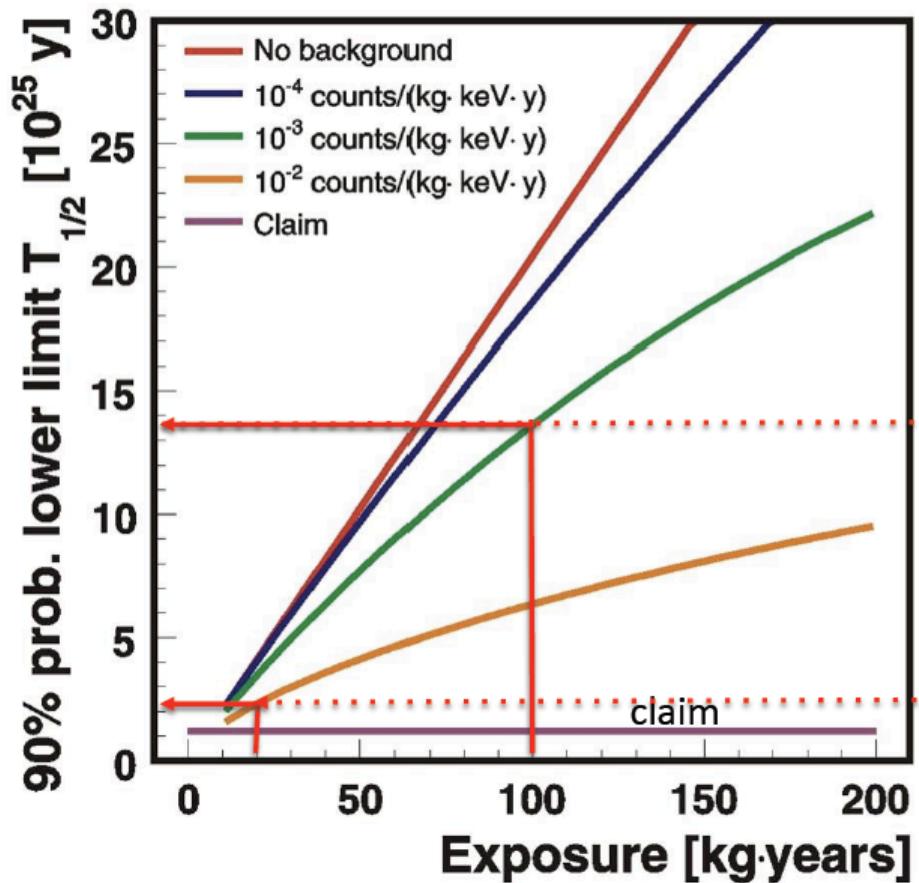


Figure adapted from:
Stefan Schönert
(for the GERDA collaboration)
LNGS Seminar, July 16, 2013

GERDA phase I background

GERDA: "background index ranging from 17.6 to 23.8×10^{-3} cts/(keV*kg*yr)", arXiv:1306.5084

Final results with PSD: background index ranging from 5 to 30×10^{-3} cts/(keV*kg*yr), arXiv:1307.4720



Phase II:

Add new enr. BEGe detectors (20 kg)
 $BI \approx 0.001$ cts / (keV kg yr) i.e. 1ct/keV t yr
Sensitivity after 100 kg yr

Phase I:

Use refurbished HdM & IGEX (18 kg)
 $BI \approx 0.01$ cts / (keV kg yr) i.e. 10 cts/keV t yr
Sensitivity after 20 kg yr

Figure adapted from:
Stefan Schönert
(for the GERDA collaboration)
LNGS Seminar, July 16, 2013

GERDA phase I results

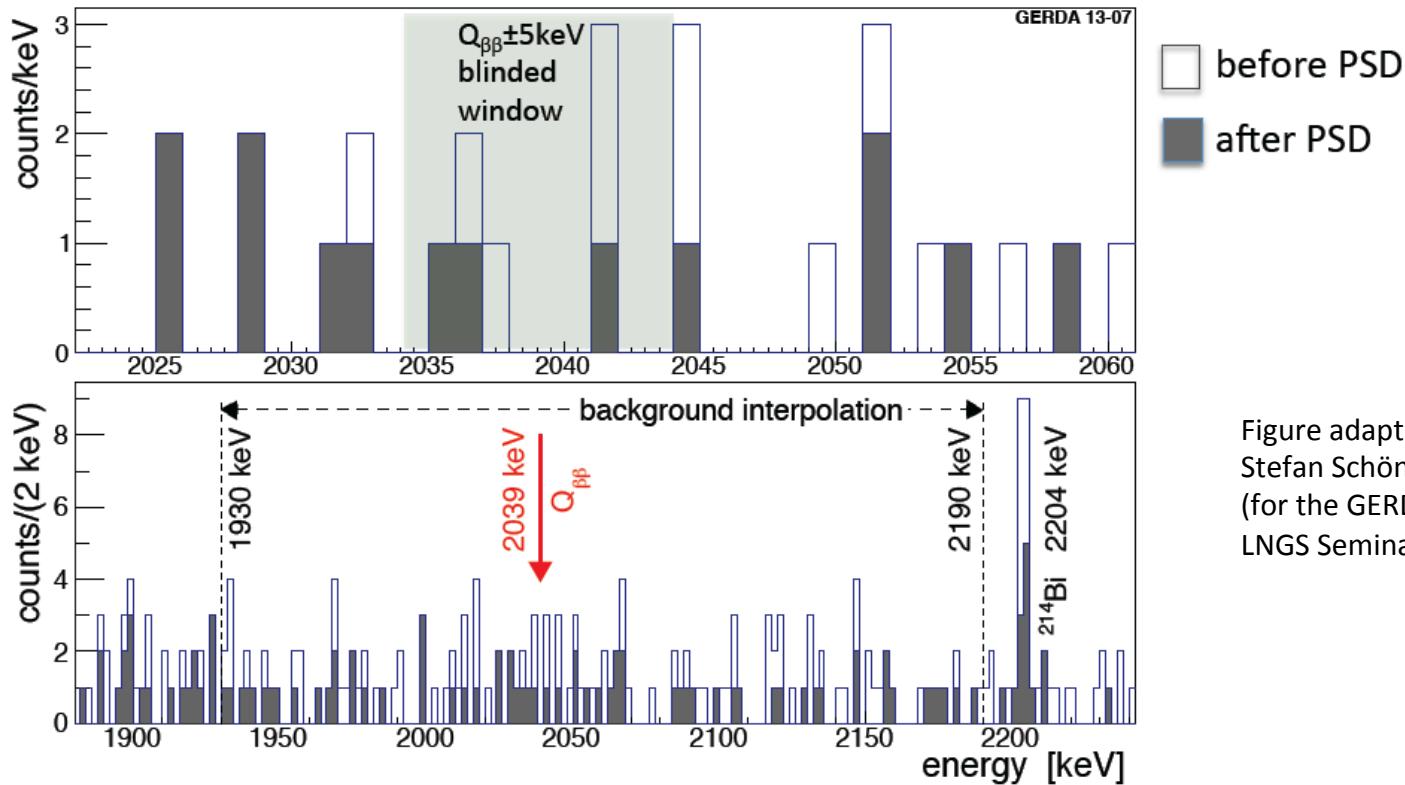


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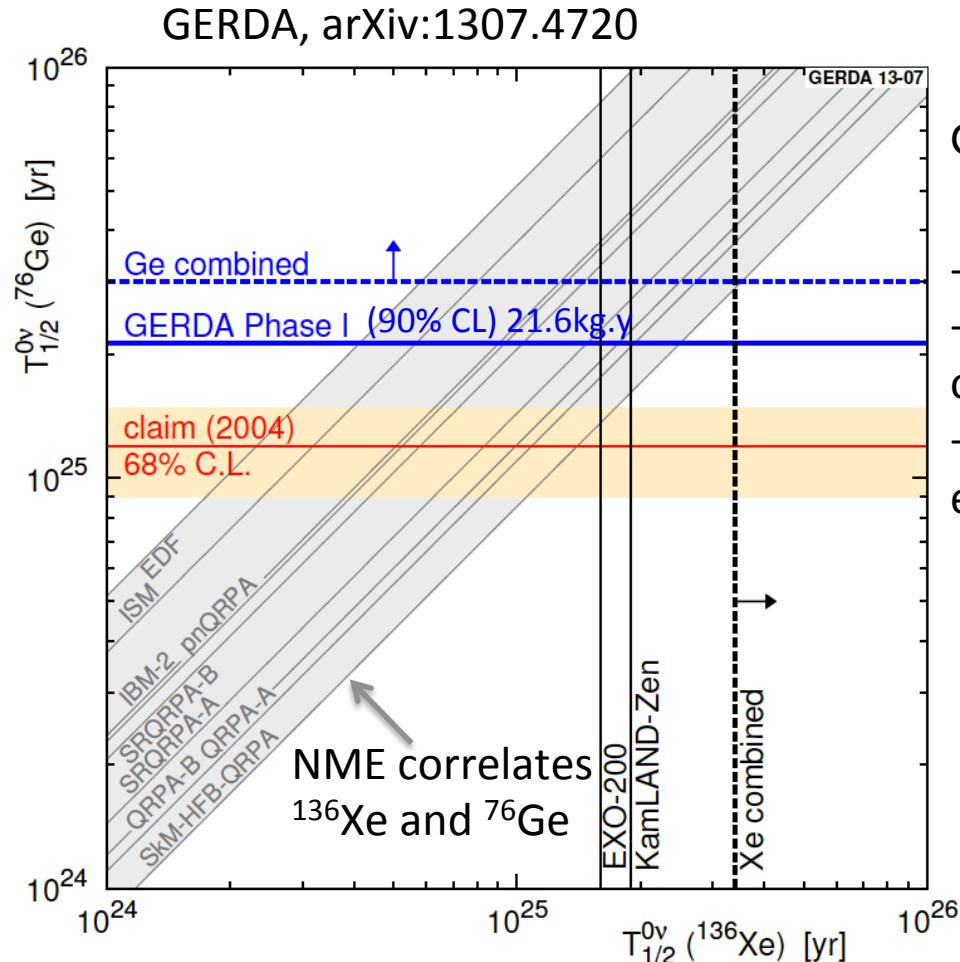
Full data set:

7 event in blinded window
 3 event survive PSD cut

Expect 5.1 background
 Expect 2.5 background

No $0\nu\beta\beta$

GERDA phase I results

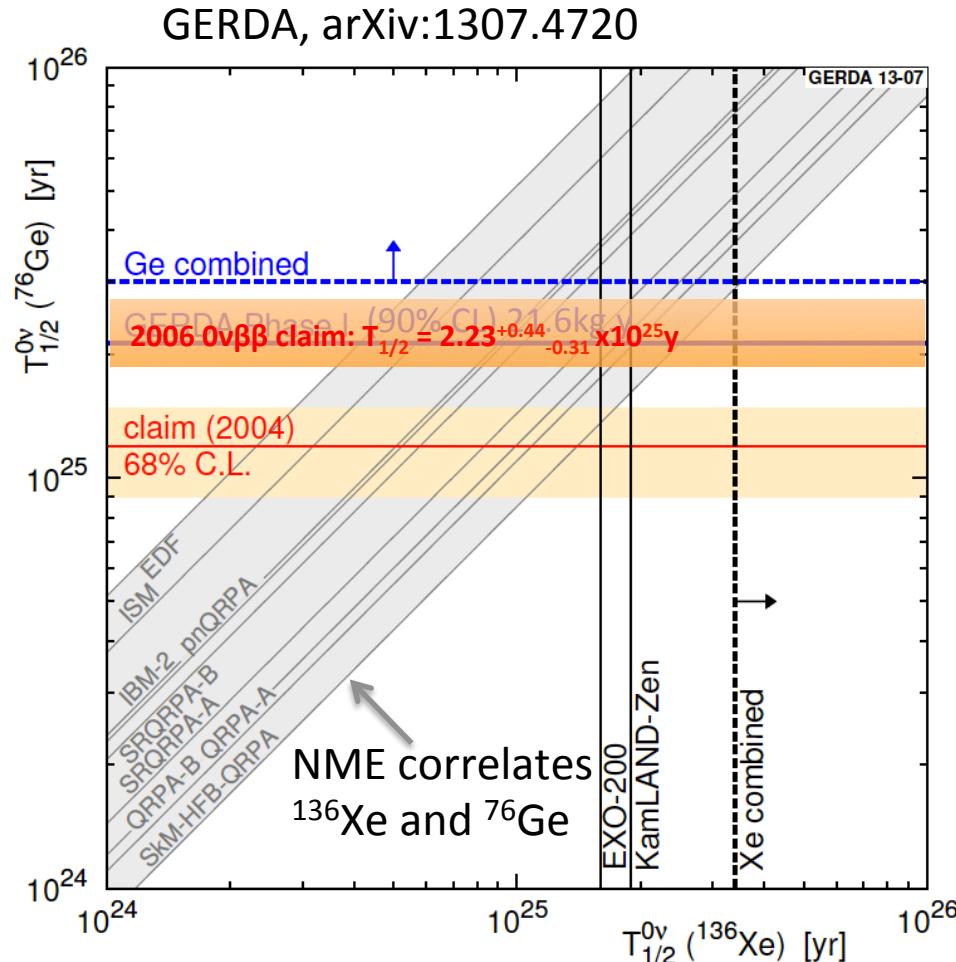


GERDA didn't compare with the 2006 claim due to "inconsistencies in the latter reference"

- The fit error on the signal count is too small.
- Expected background doesn't match observed (fluctuation probability $\sim 5 \cdot 10^{-7}$)
- Inconsistency with the signal detection efficiency

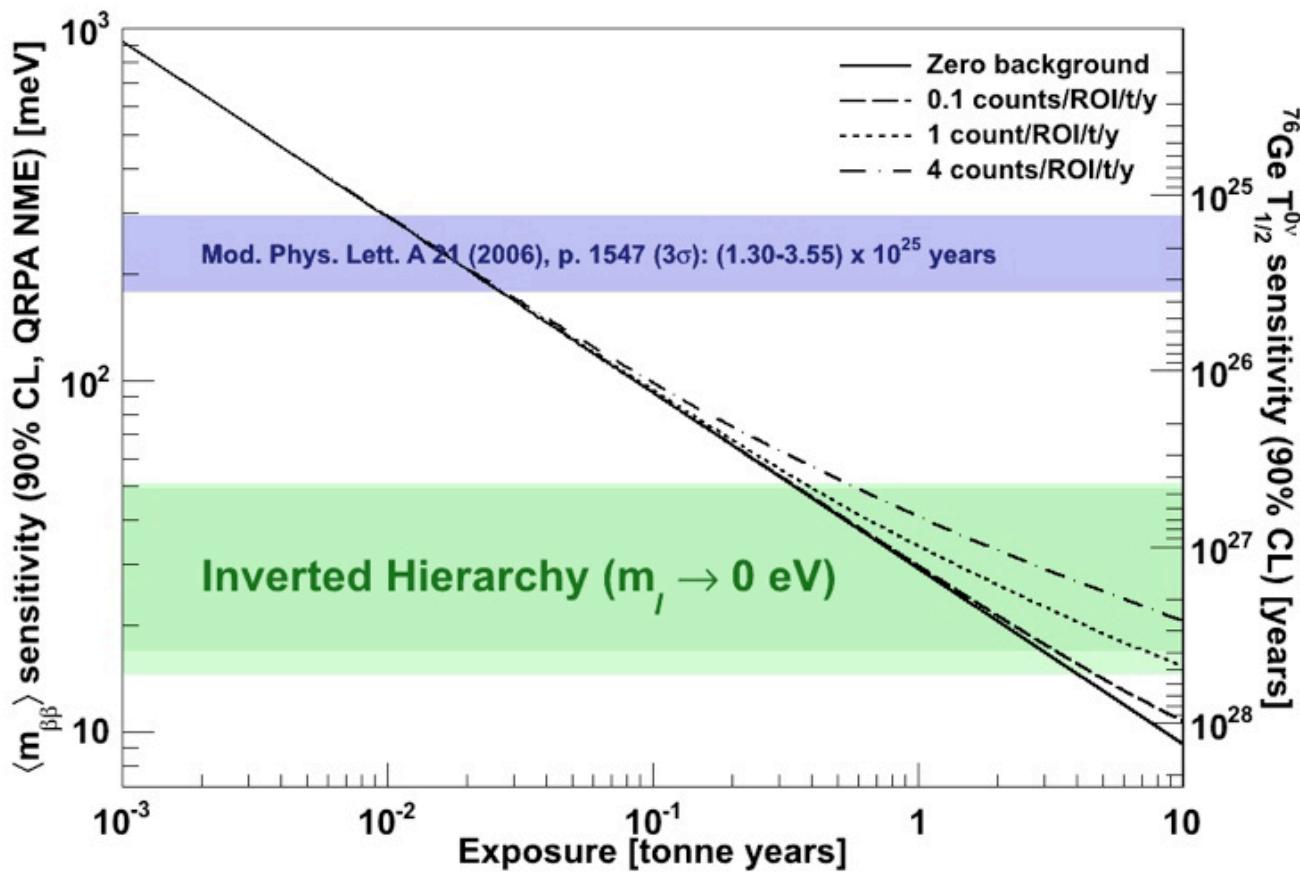
B. Schwingenheuer, Ann. Phys. (Berlin) 525, No. 4 (2013)

GERDA phase I results

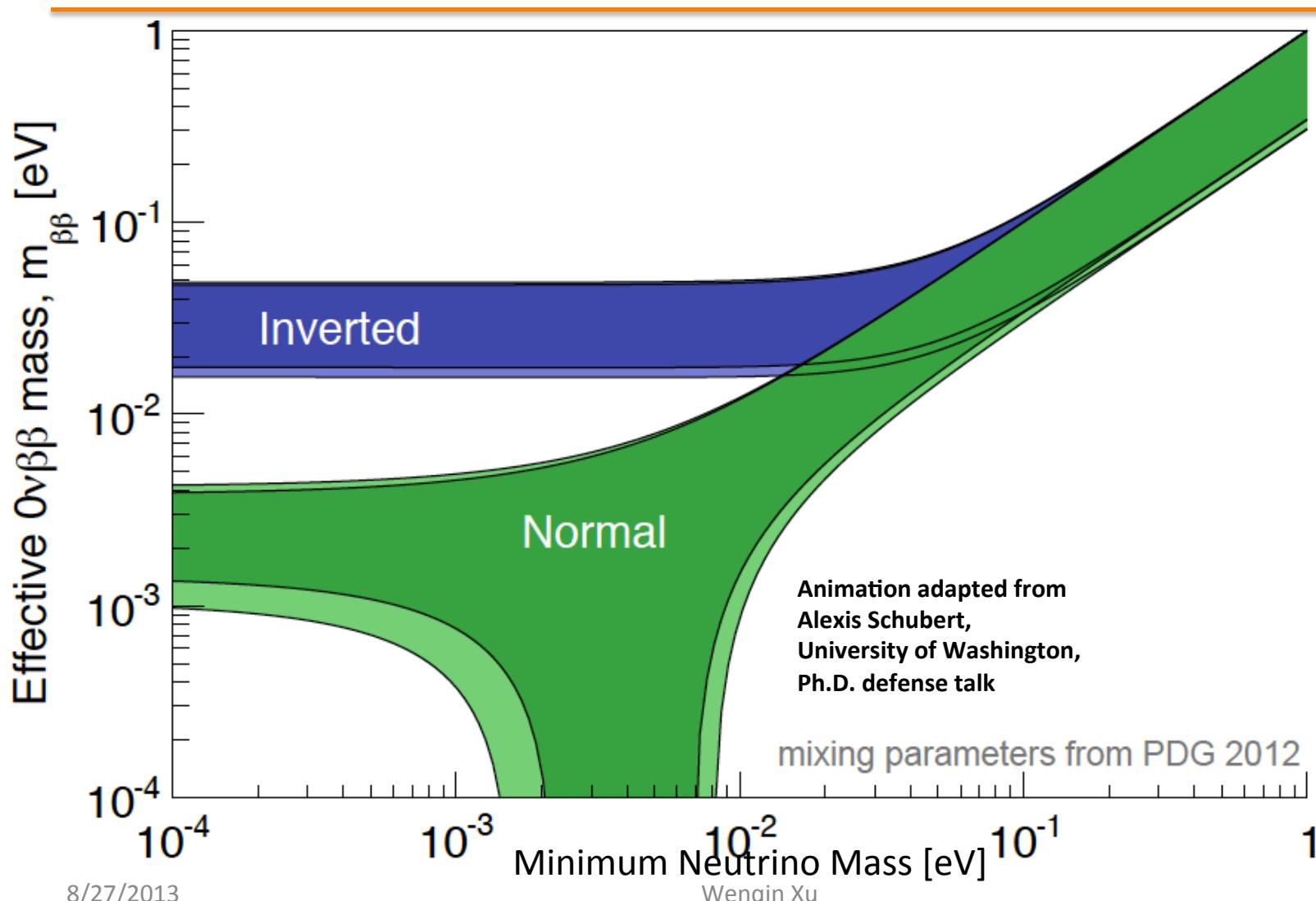


MAJORANA DEMONSTRATOR

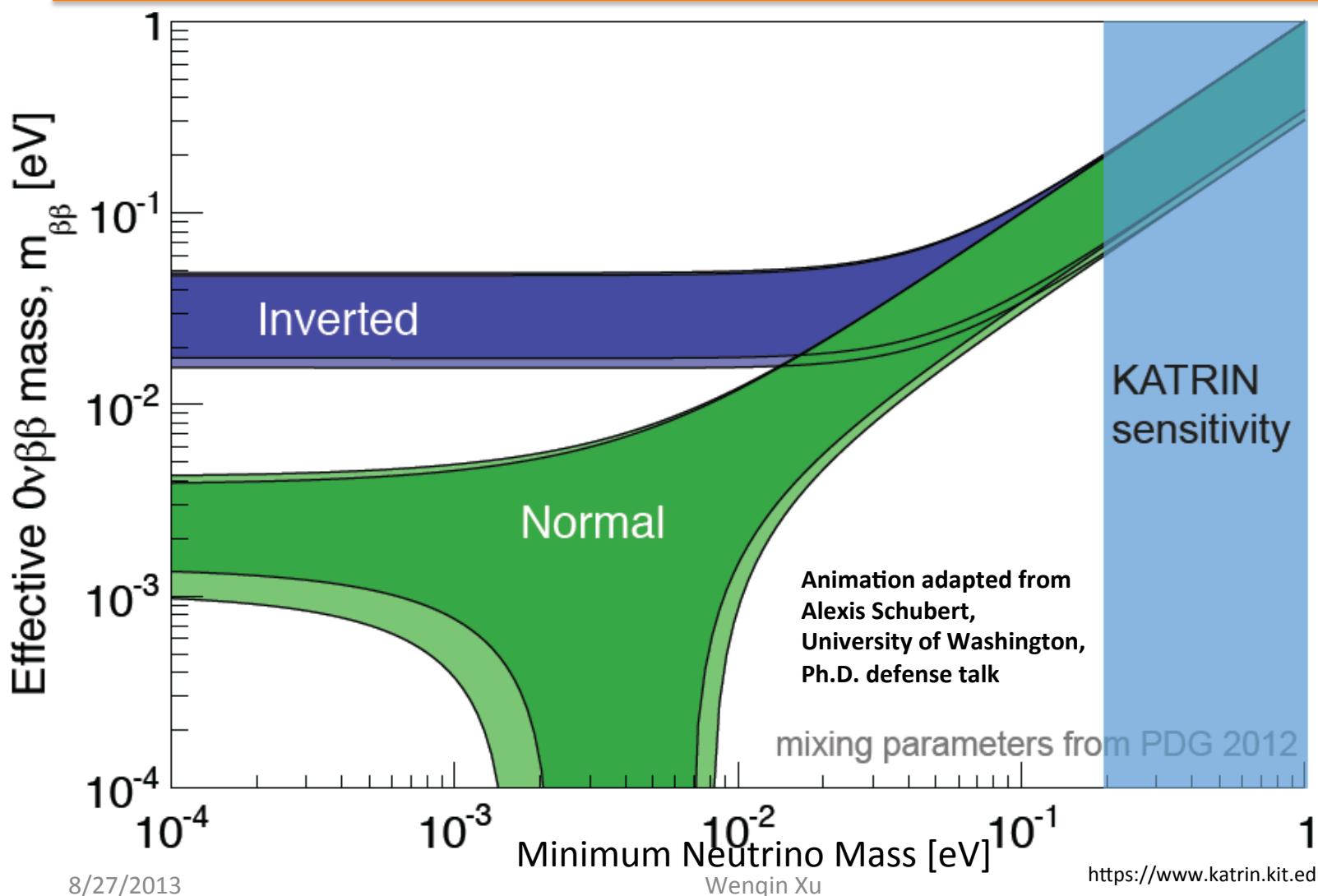
sensitivity



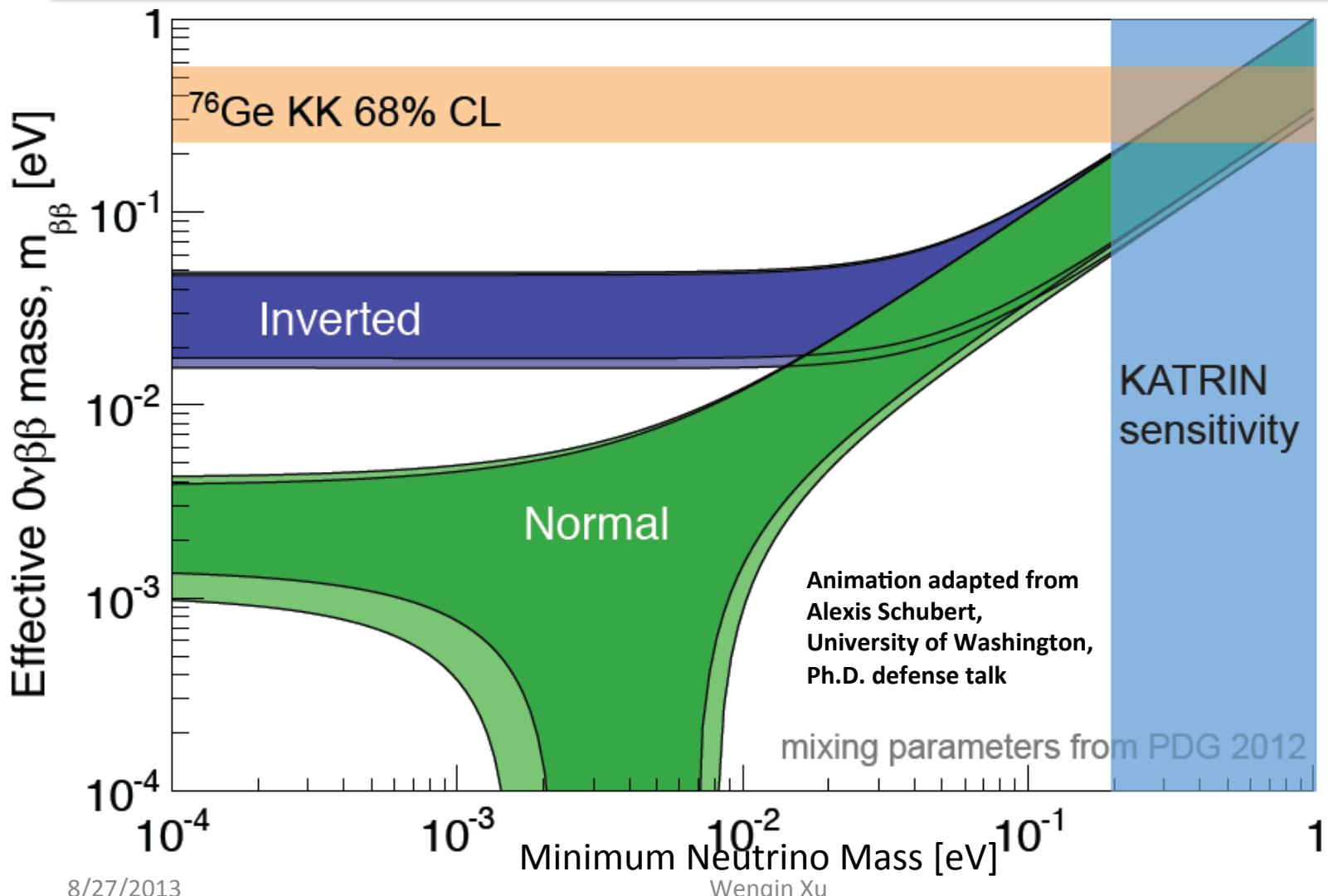
The mass plot



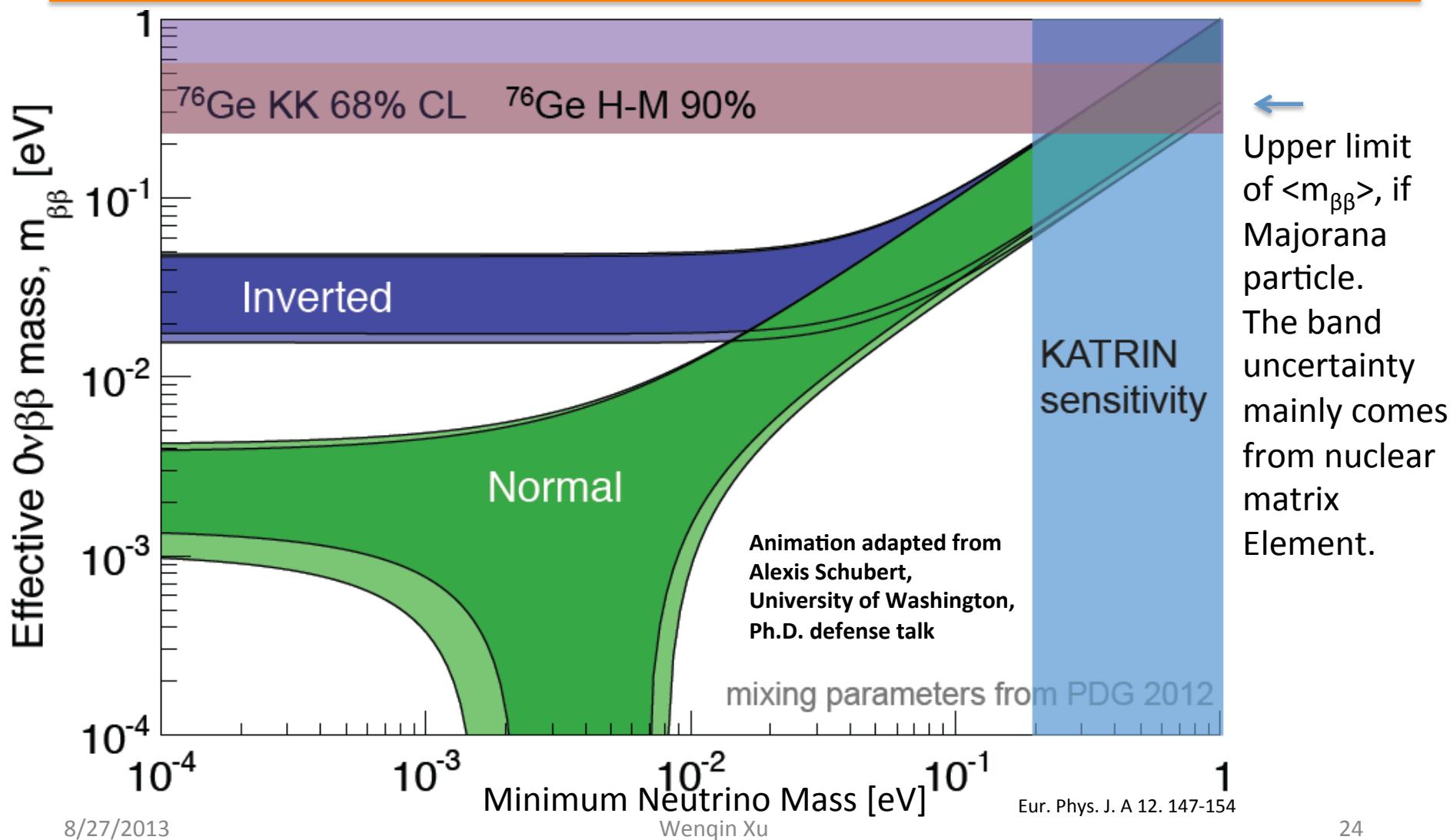
KATRIN sensitivity



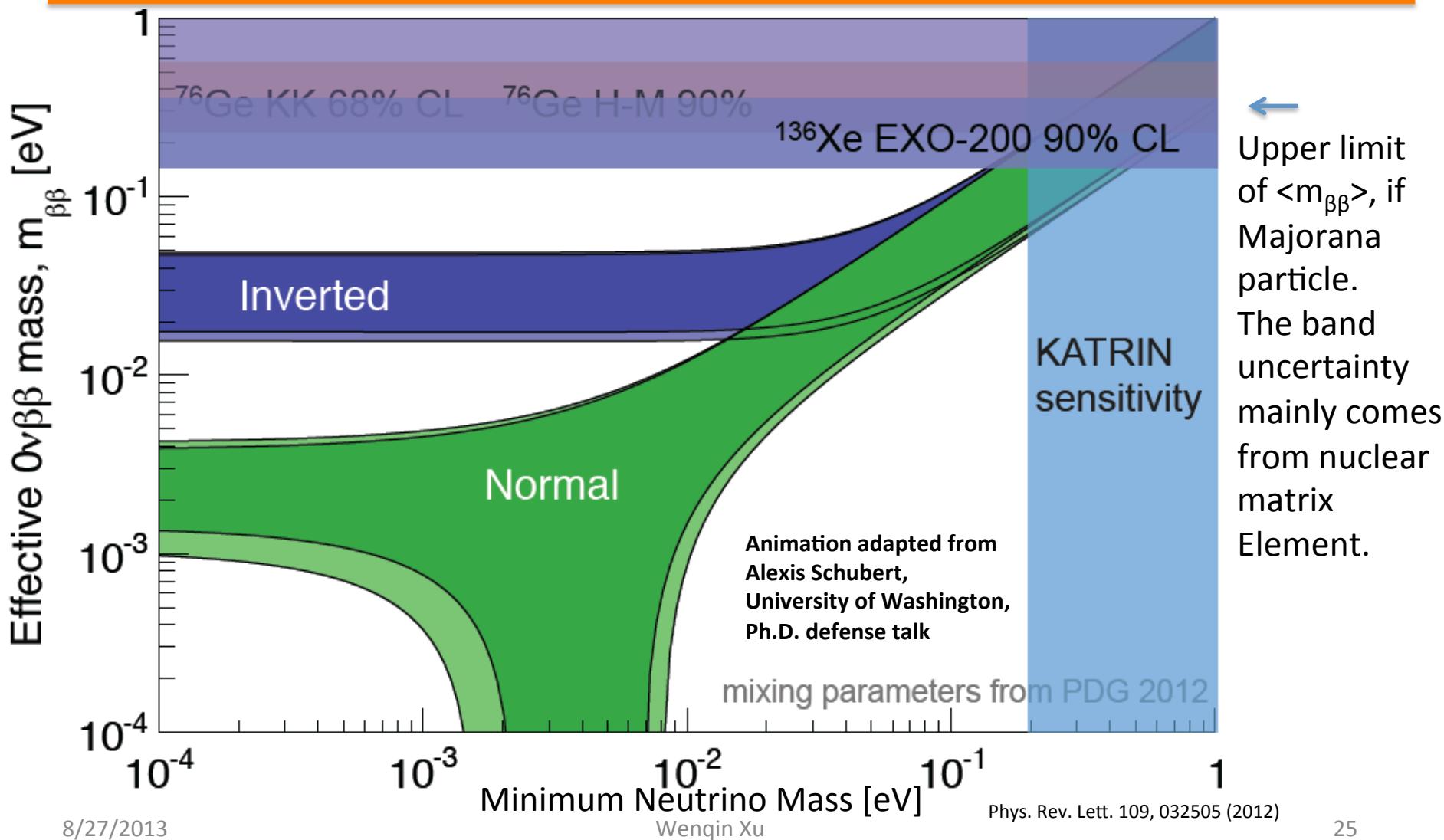
Neutrinoless double beta decay



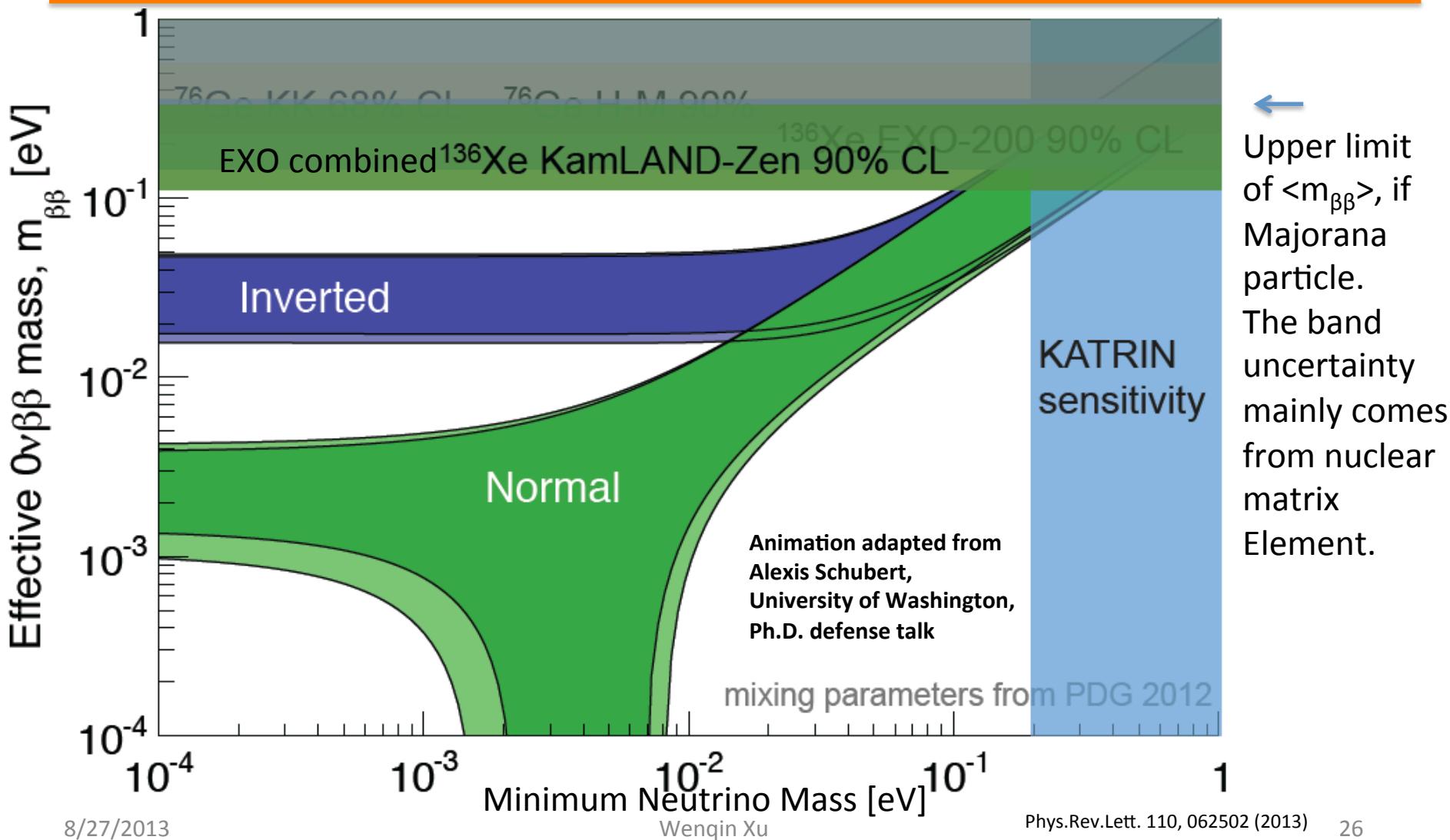
Neutrinoless double beta decay



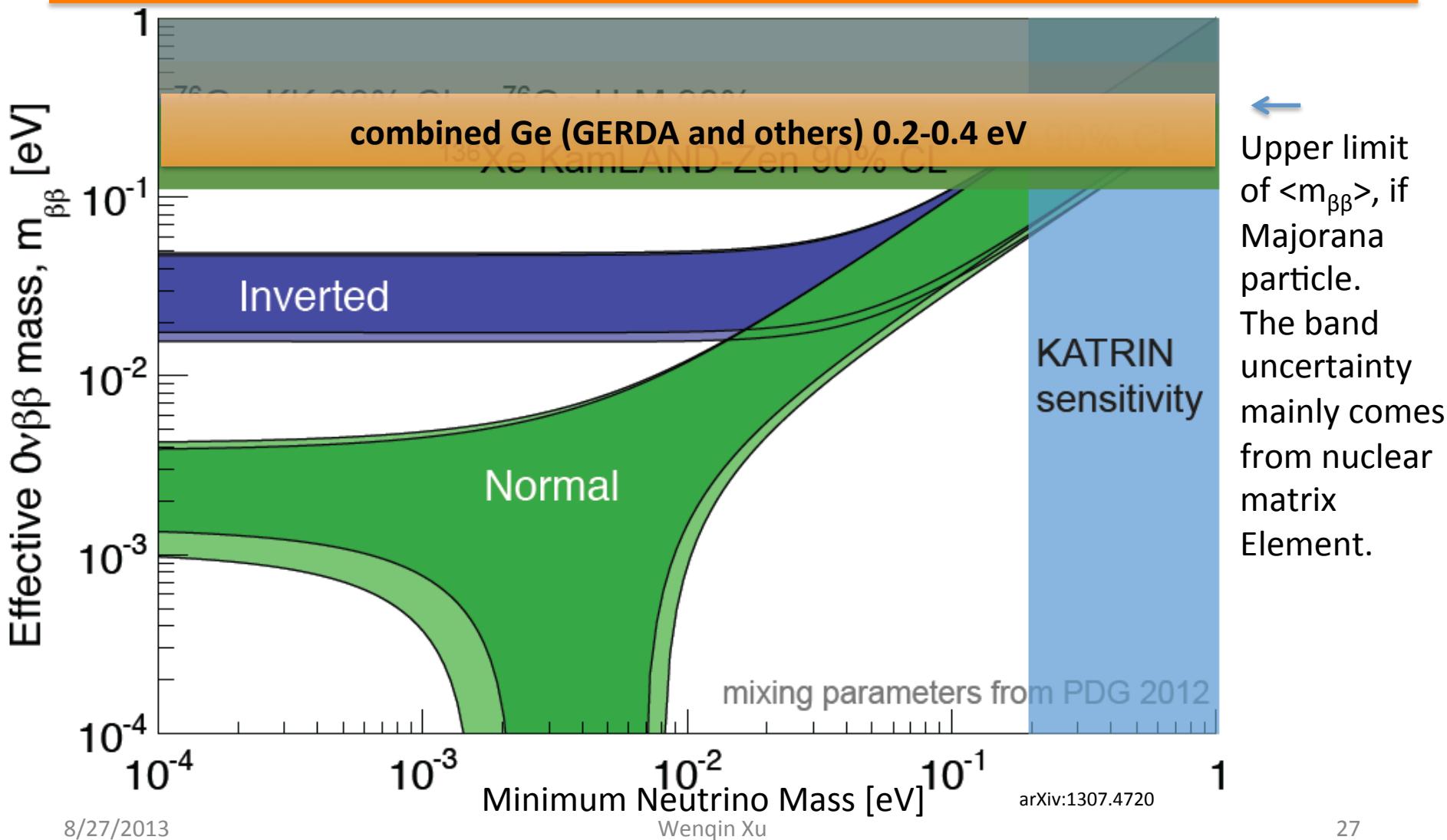
Neutrinoless double beta decay



Neutrinoless double beta decay

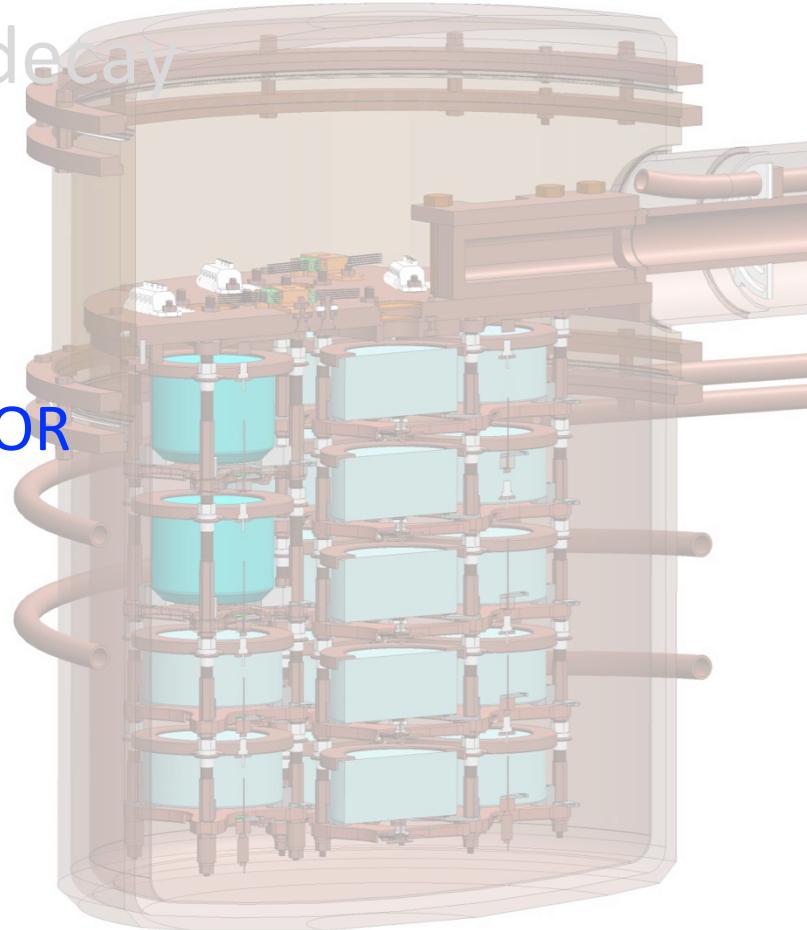


Neutrinoless double beta decay



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 - the physics
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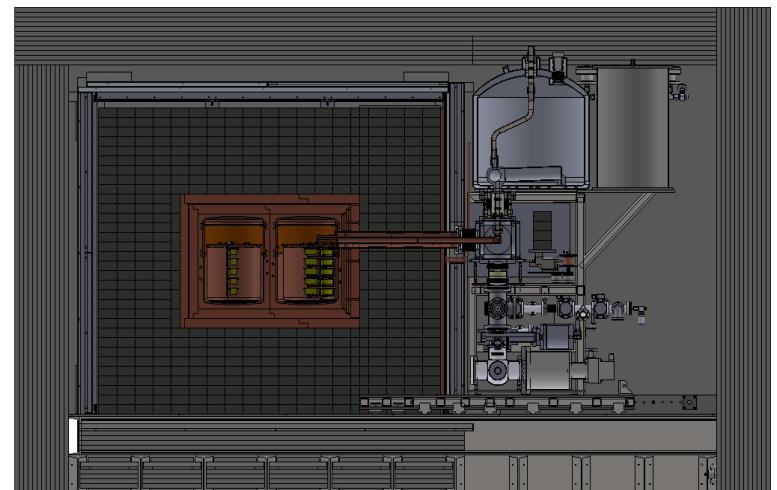
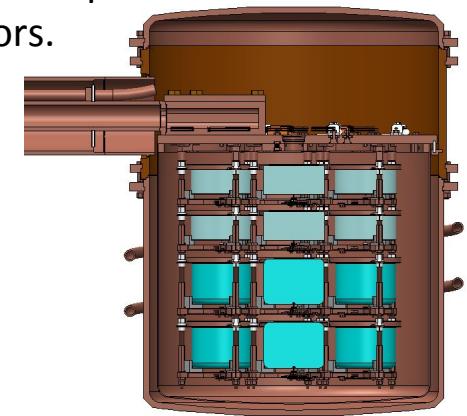


The MAJORANA DEMONSTRATOR

Funded by DOE Office of Nuclear Physics and NSF Particle Astrophysics,
with additional contributions from international collaborators.

- Goals:**
- Demonstrate backgrounds low enough to justify building a tonne scale experiment.
 - Establish feasibility to construct & field modular arrays of Ge detectors.
 - Test Klapdor-Kleingrothaus claim.
 - Low-energy dark matter (light WIMPs, axions, ...) searches.

- Located underground at 4850' Sanford Underground Research Facility
- Background Goal in the $0\nu\beta\beta$ peak region of interest (4 keV at 2039 keV)
[3 counts/ROI/t/y \(after analysis cuts\)](#)
scales to 1 count/ROI/t/y for a tonne experiment
- 40-kg of Ge detectors
 - 30 kg of 86% enriched ^{76}Ge crystals &
10 kg of $^{\text{nat}}\text{Ge}$
 - Detector Technology: P-type, point-contact.
- 2 independent cryostats
 - ultra-clean, electroformed Cu
 - 20 kg of detectors per cryostat
 - naturally scalable
- Compact Shield
 - low-background passive Cu and Pb
shield with active muon veto

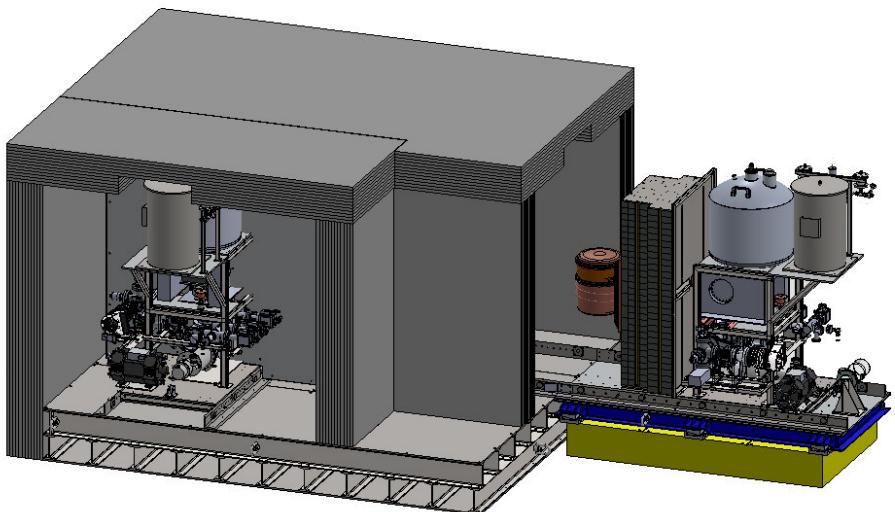
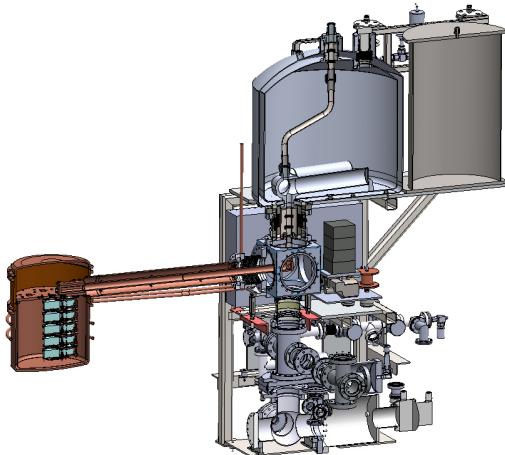




MJD Implementation

- Three Steps

- Prototype Cryostat* (2 strings, ^{nat}Ge)
- Cryostat 1 (3 strings ^{enr}Ge & 4 strings ^{nat}Ge)
- Cryostat 2 (7 strings ^{enr}Ge)



* Same design as Cryos 1 & 2, but fabricated using OFHC Cu (non-electroformed) components.

Sanford Underground Research Facility (SURF), Lead, SD



<http://sanfordlab.org>

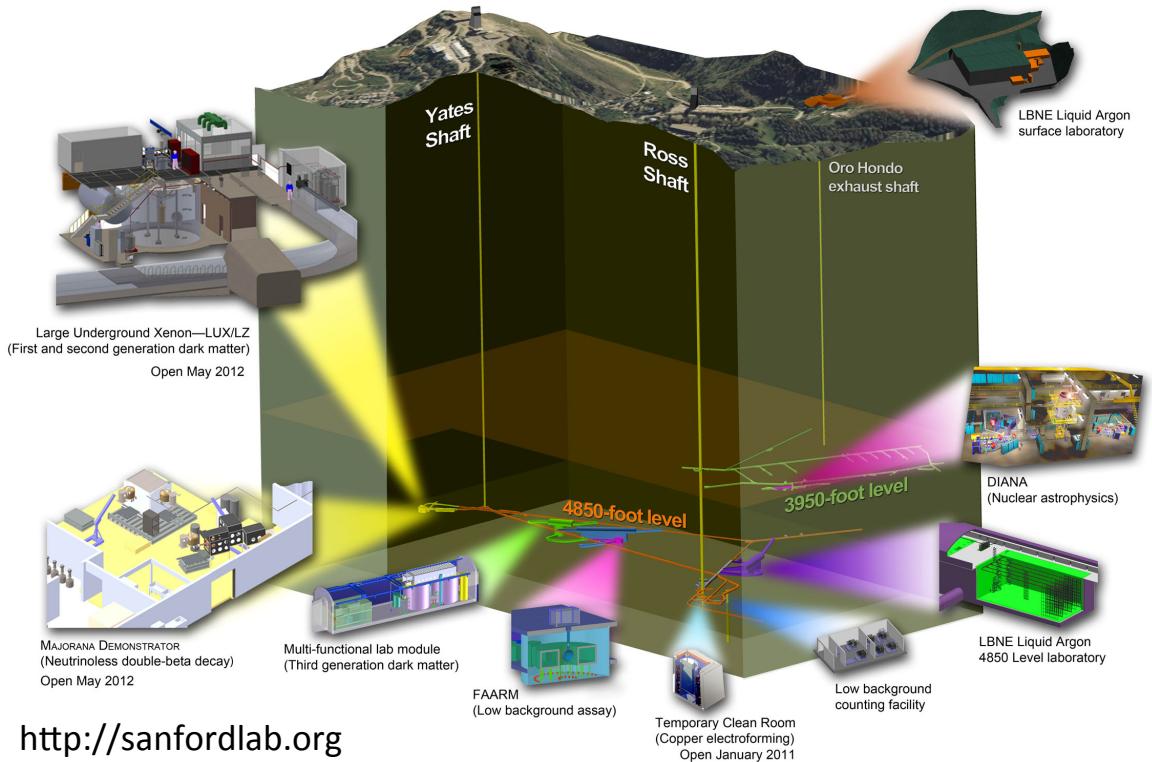


Sanford Underground Research Facility (SURF), Lead, SD

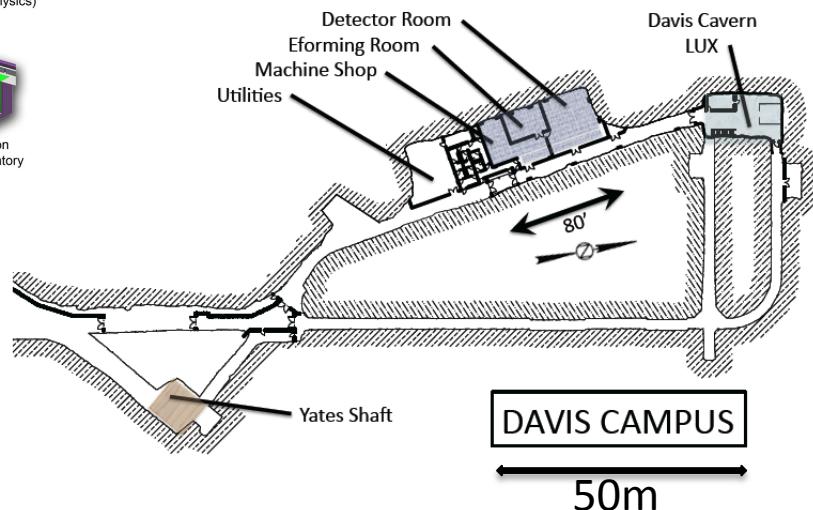


Photo Courtesy of R. Martin

Sanford Underground Research Facility (SURF), Lead, SD



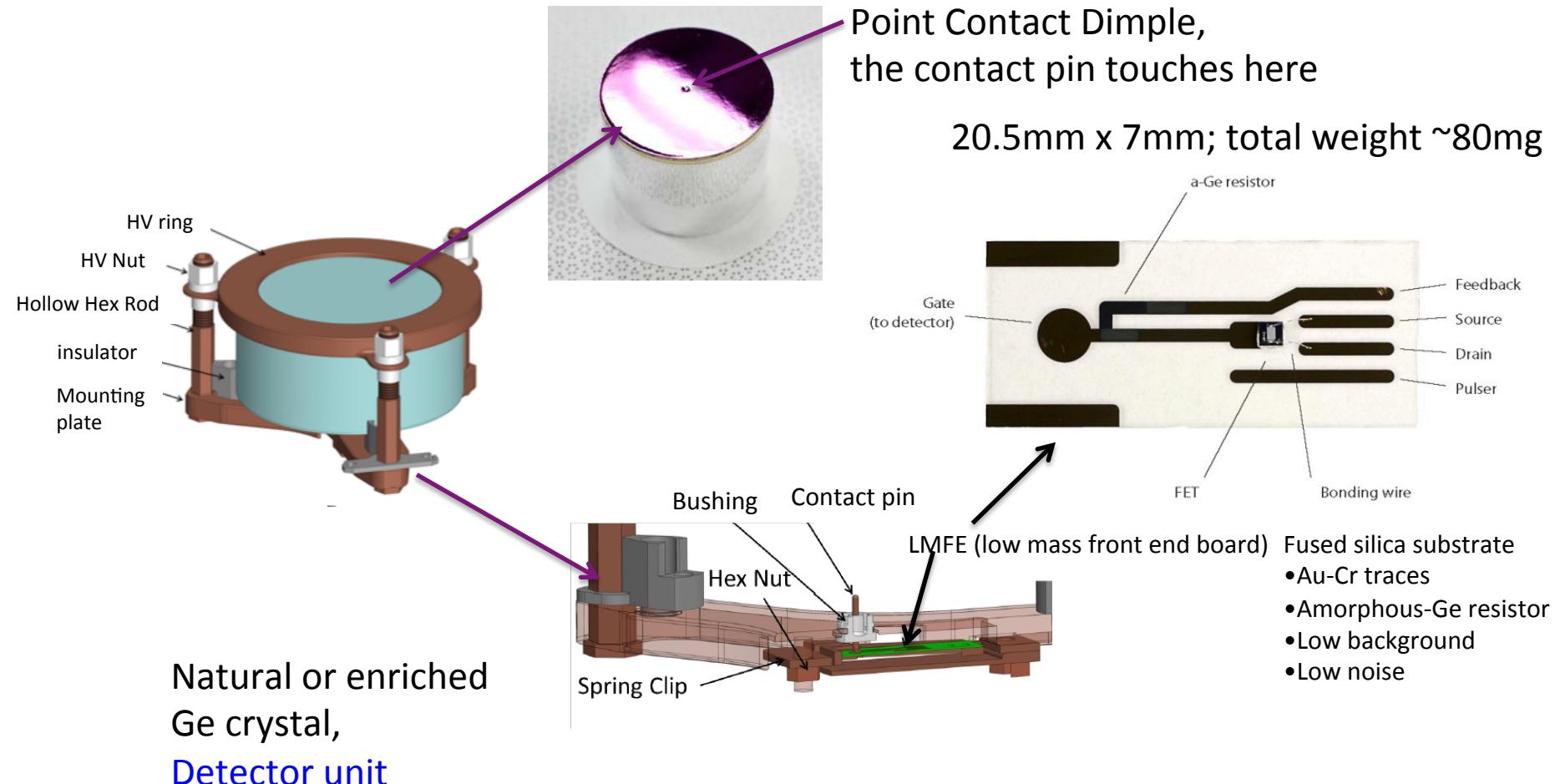
4850 foot ~ 4260 m.w.e



<http://sanfordlab.org>

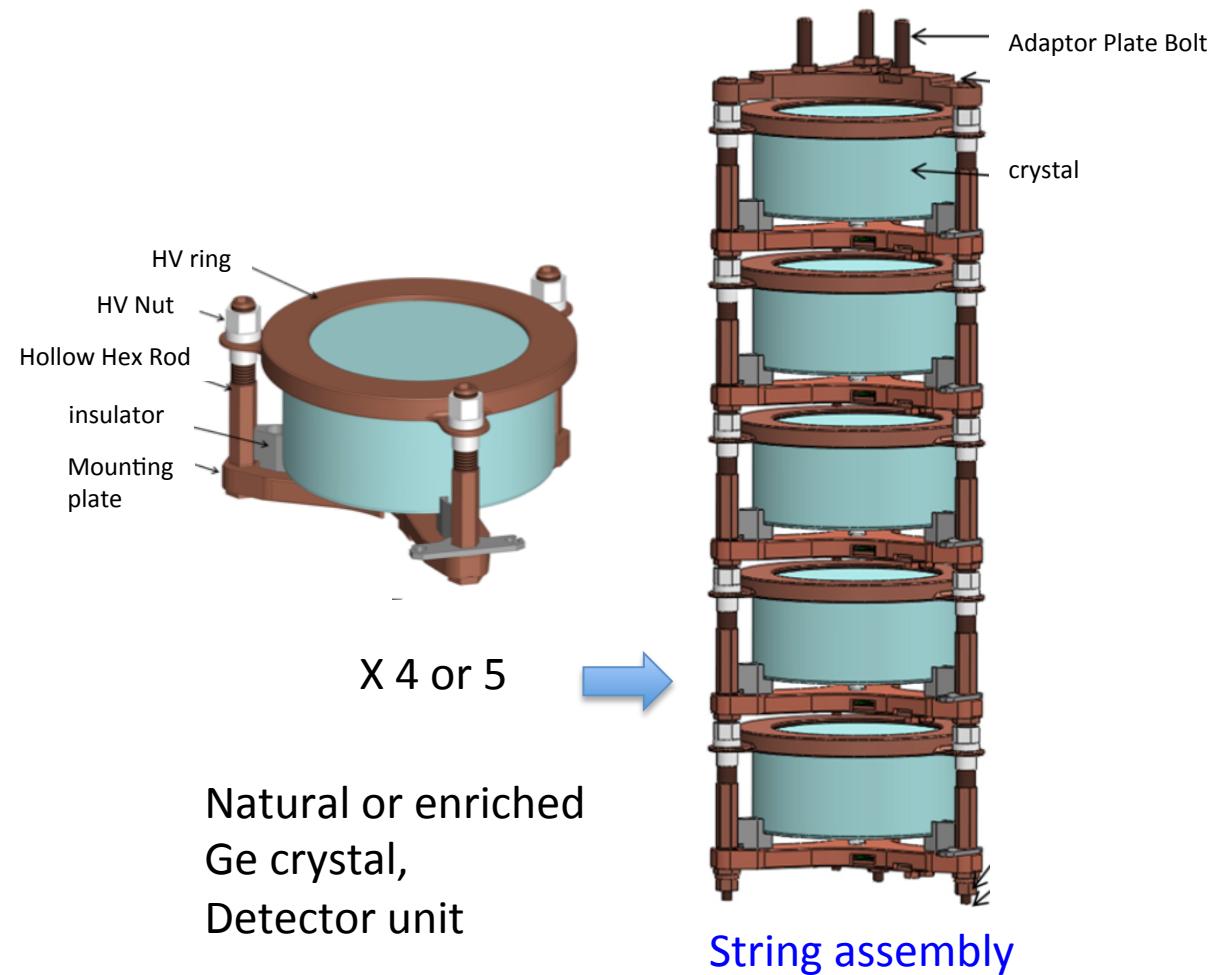


Detector Unit Assemble



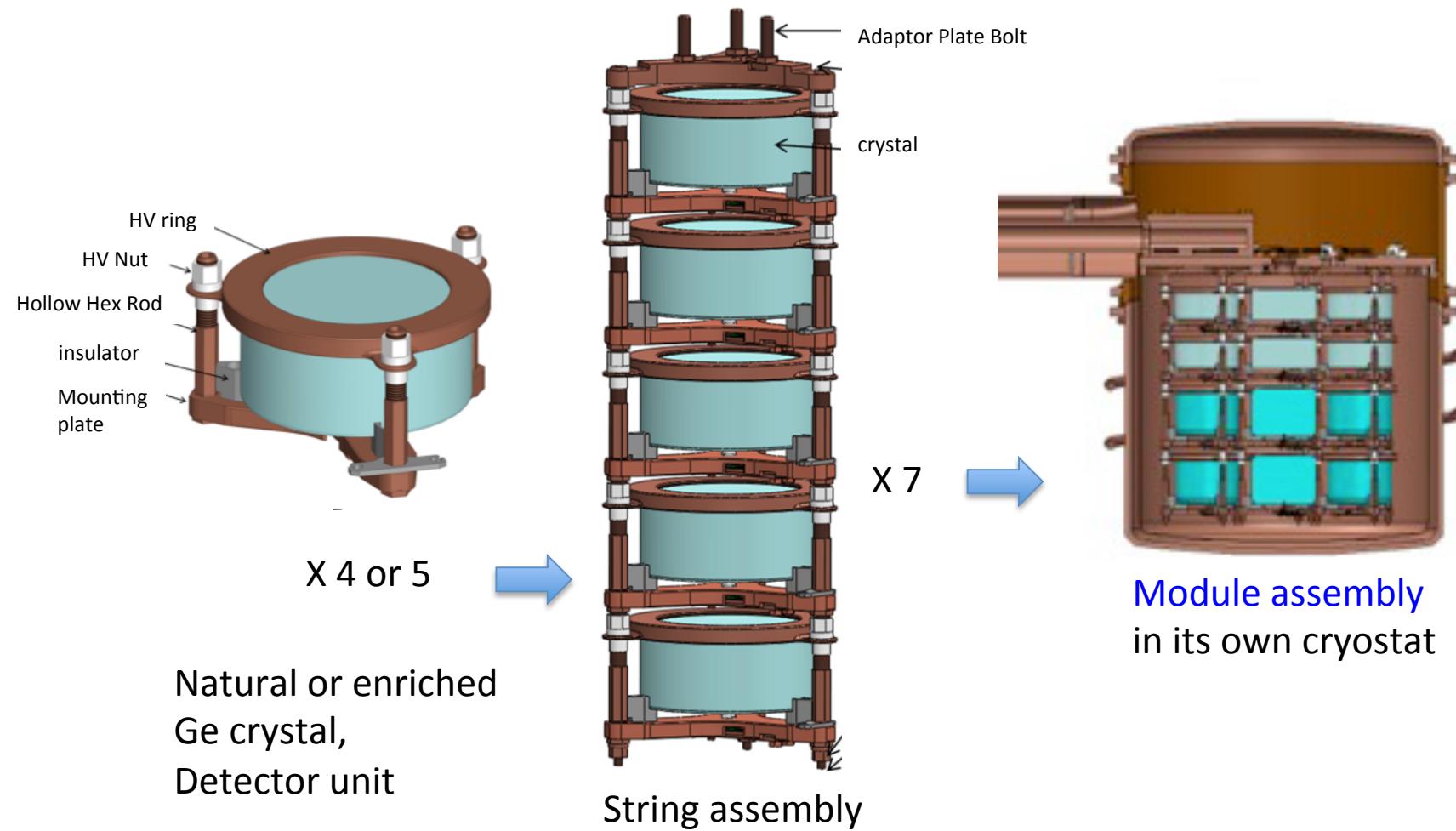


Modular approach



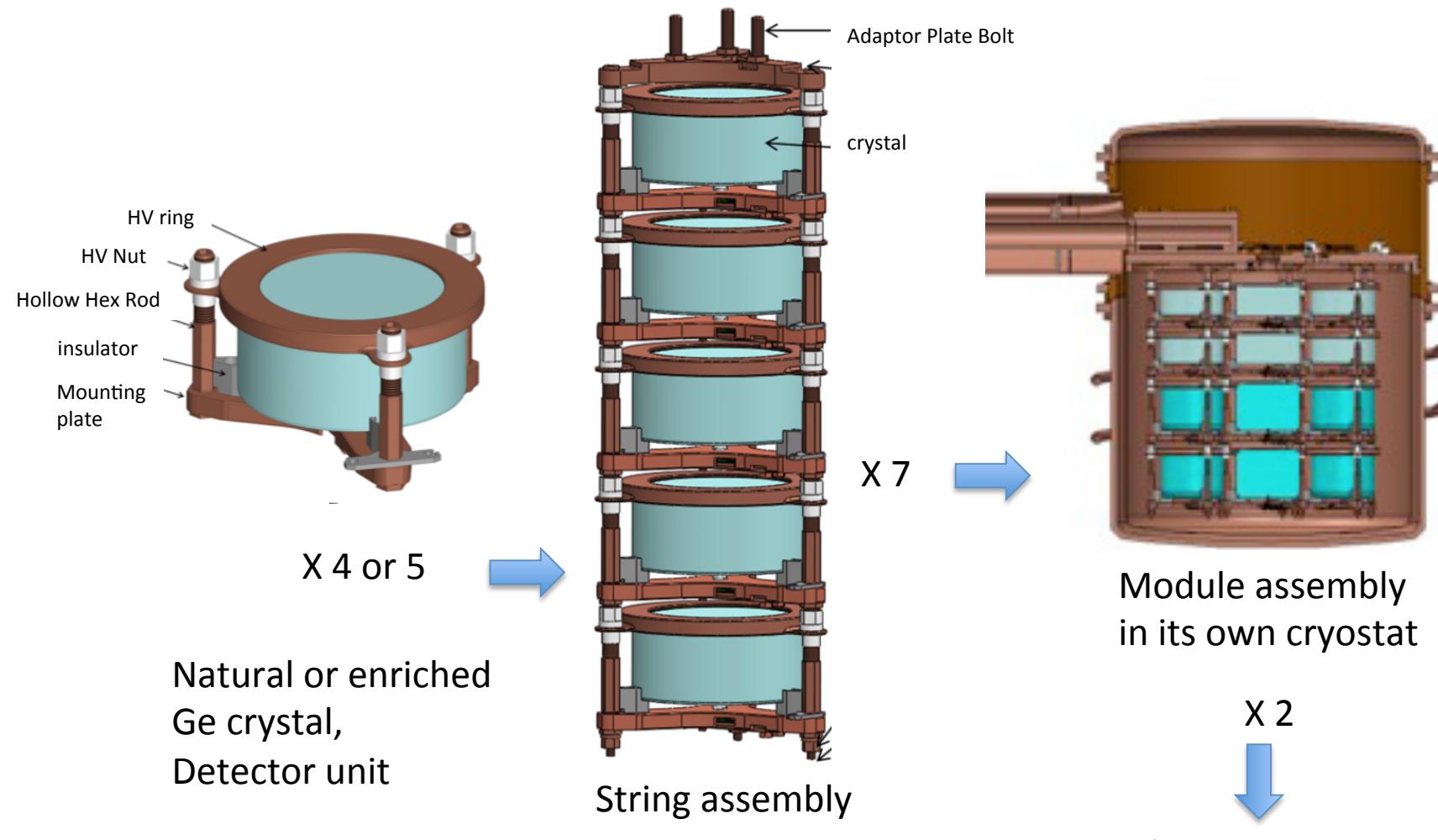


Modular approach



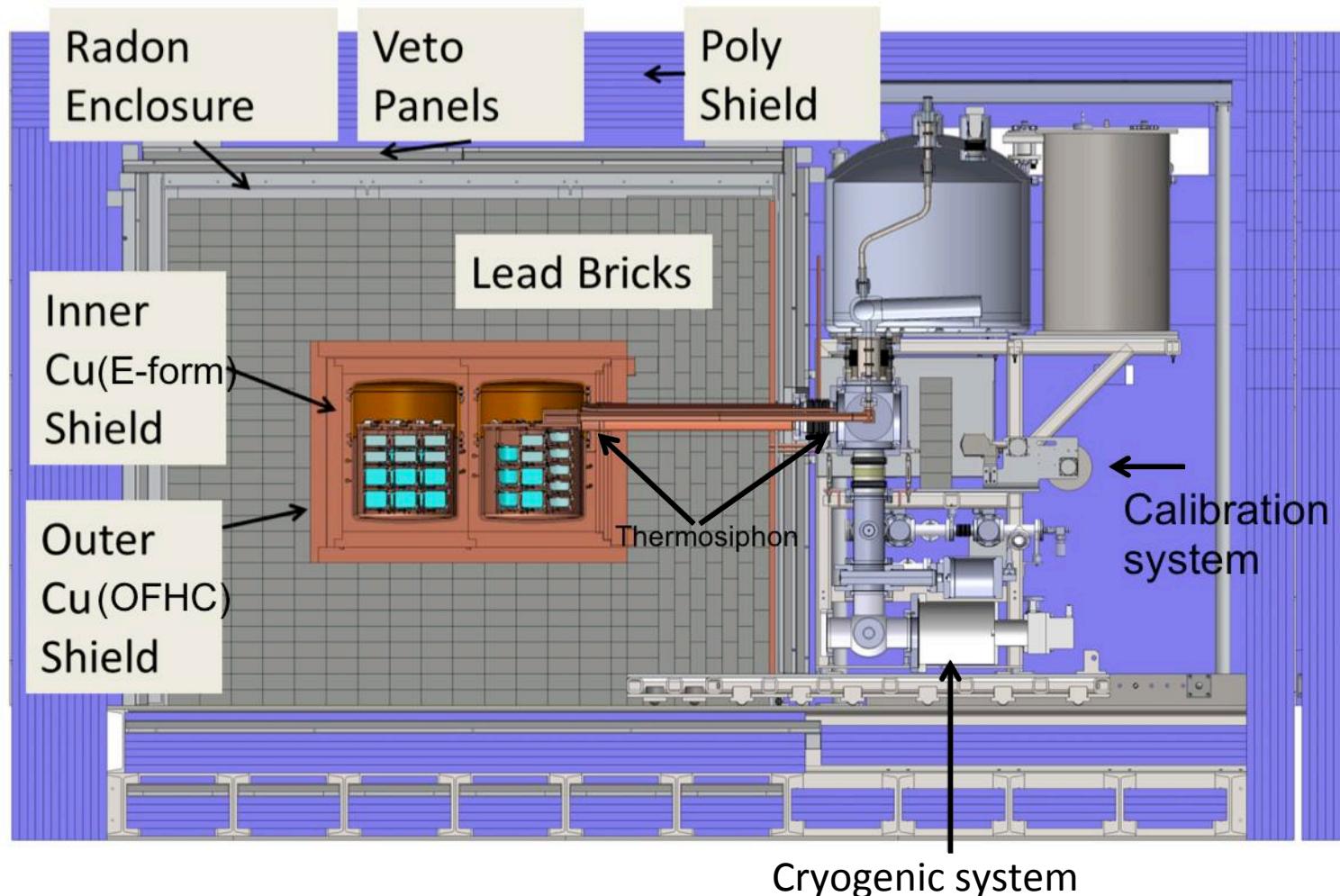


Modular approach





The shielding



The veto panels (scintillating acrylic 2 layers, 2.54cm): **active** shielding.

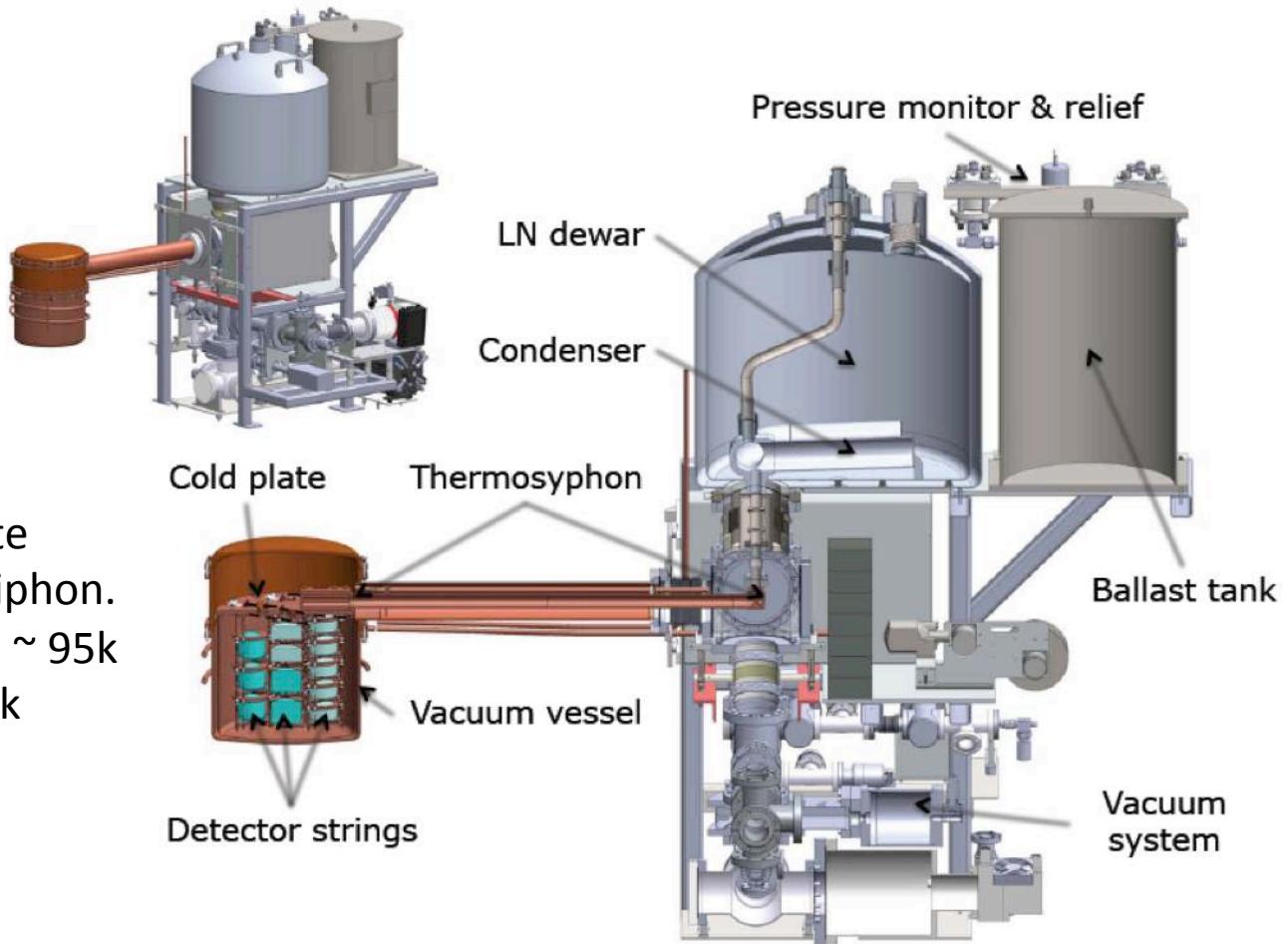
The others are **passive**:

- Poly shield - 30cm;
- Radon enclosure - 0.32-0.635cm Al (purged with N₂);
- Lead - 45 cm;
- Outer Cu - 5 cm;
- Inner Cu 4 layers - 1.25 cm each.

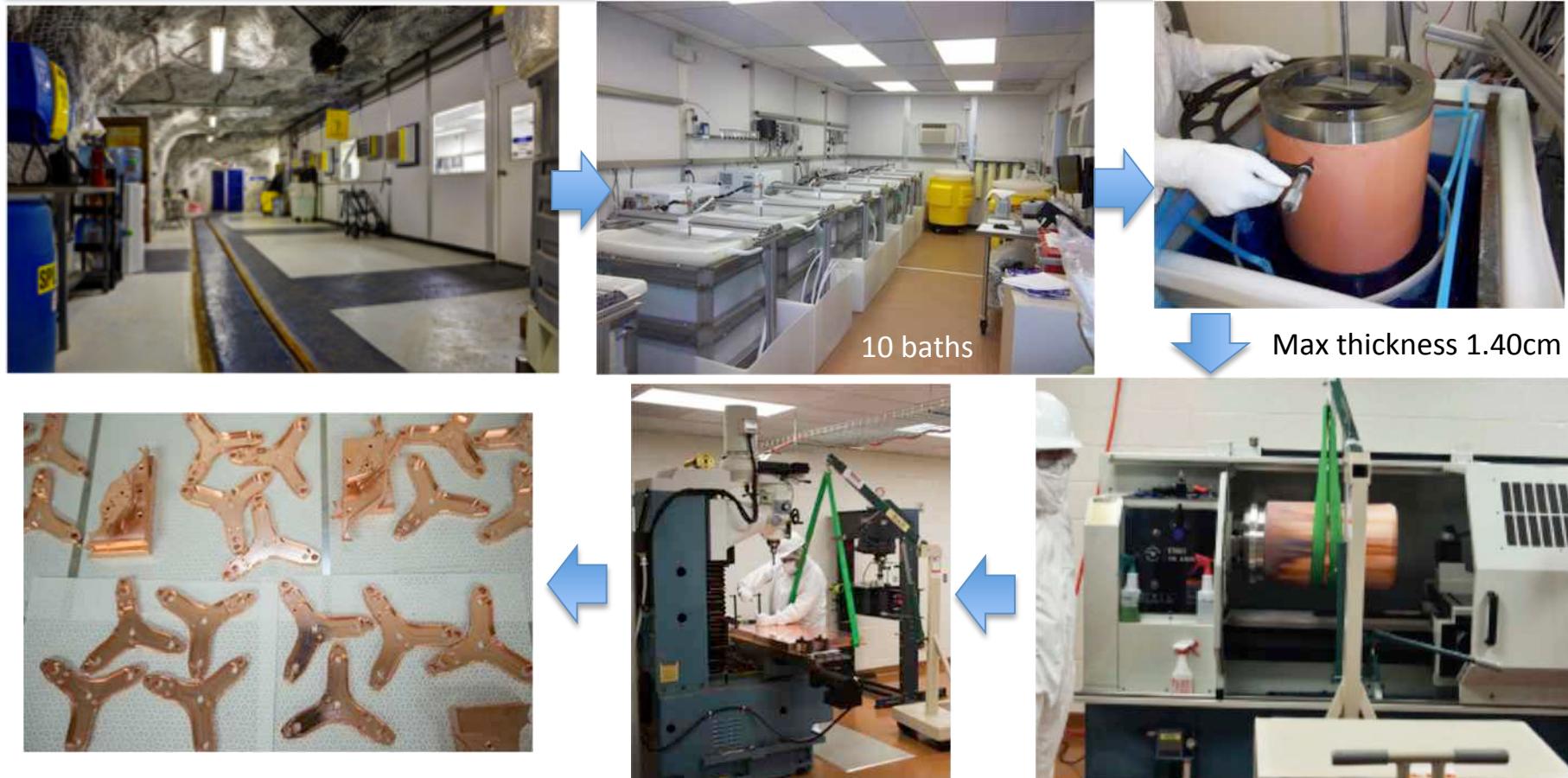


Cryogenic System

Cooling to the cold plate provided by a thermosiphon.
Detectors temperature $\sim 95\text{K}$
FET temperature $\sim 150\text{K}$
(self-heating)



Electroformed copper (EFCu)



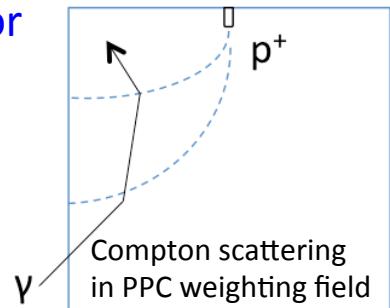
Assay results for e-form copper ^{232}Th : $0.7 \pm 0.3 \mu\text{Bq/kg}$. ^{238}U : $<1.3 \mu\text{Bq/kg}$
One order of magnitude better than the cleanest commercial copper



Point Contact Detector

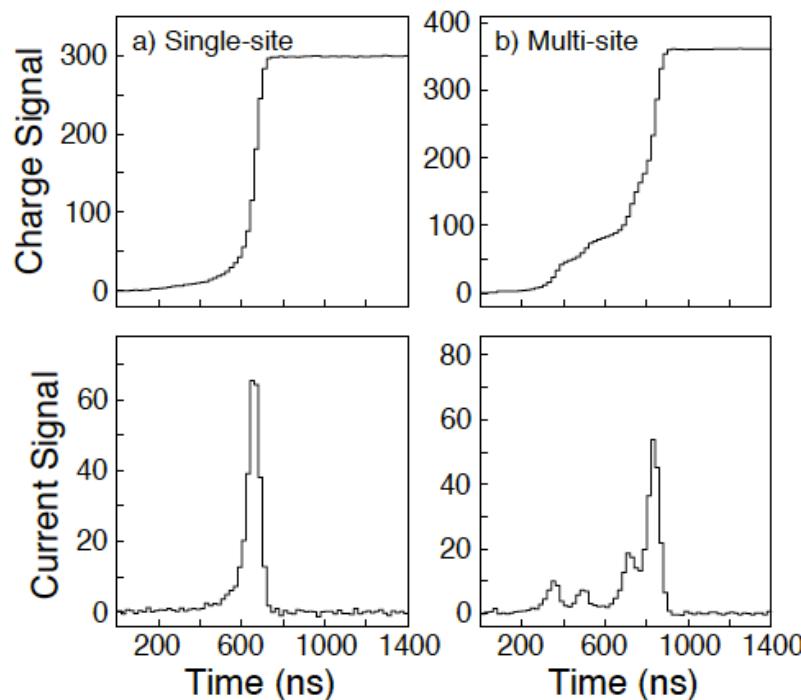


P-type Point Contact Detector
(PPC)



Point Contact Detector:

- Relatively low electrical fields
- Larger time spreads for spatially distinct energy depositions
- Crucial in distinguishing multi-site γ background from single-site β signals



Pulse-Shape-Discrimination (PSD)



Pulse-Shape-Discrimination

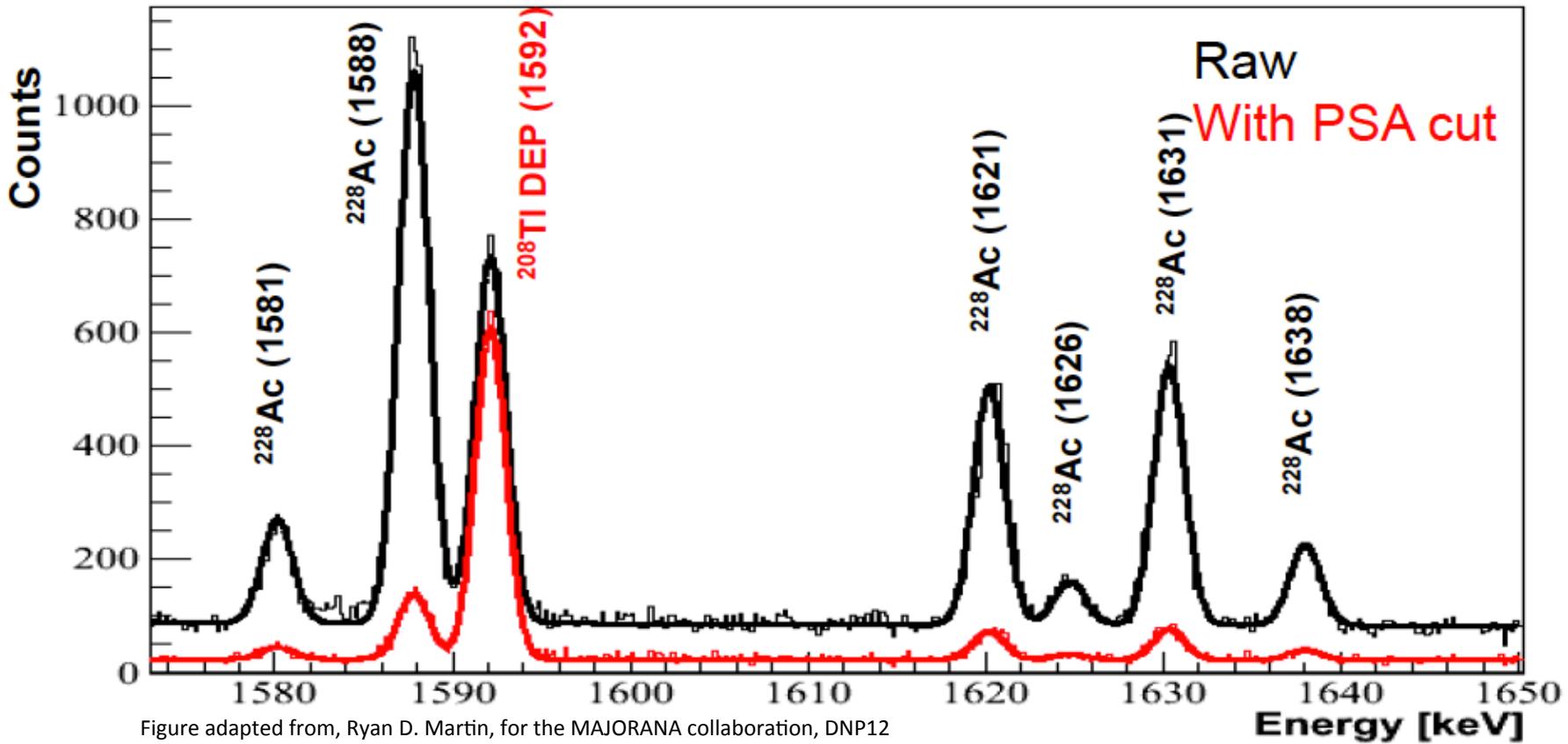


Figure adapted from, Ryan D. Martin, for the MAJORANA collaboration, DNP12

Retain 90% Double Escape Peak: single-site events, similar to $0\nu\beta\beta$ and $2\nu\beta\beta$
Reject 89% Full Energy Peaks: multi-site events, background-like



Low background is the key

Natural radioactivity:

- **Pure material** (e.g. EFCu, clean plastic and others)
- **Shielding**
- **Analysis cuts** (PSD, granularity cuts)



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Cosmogenic:

- **Deep underground** Combined efficiency of two layers of veto panel~99.9%,
Un-vetoed direct muon background<0.03 counts/ROI/t/y.
- **Muon veto**



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- **Limit surface time of Ge** (shielded shipping and storage)

Low background is the key



Delivery of enriched GeO_2
from Russia to Oak Ridge
 \sim 1 month by sea



- **Limit surface time of Ge** (shielded shipping and storage)



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- **E-form Cu underground**



Low background is the key

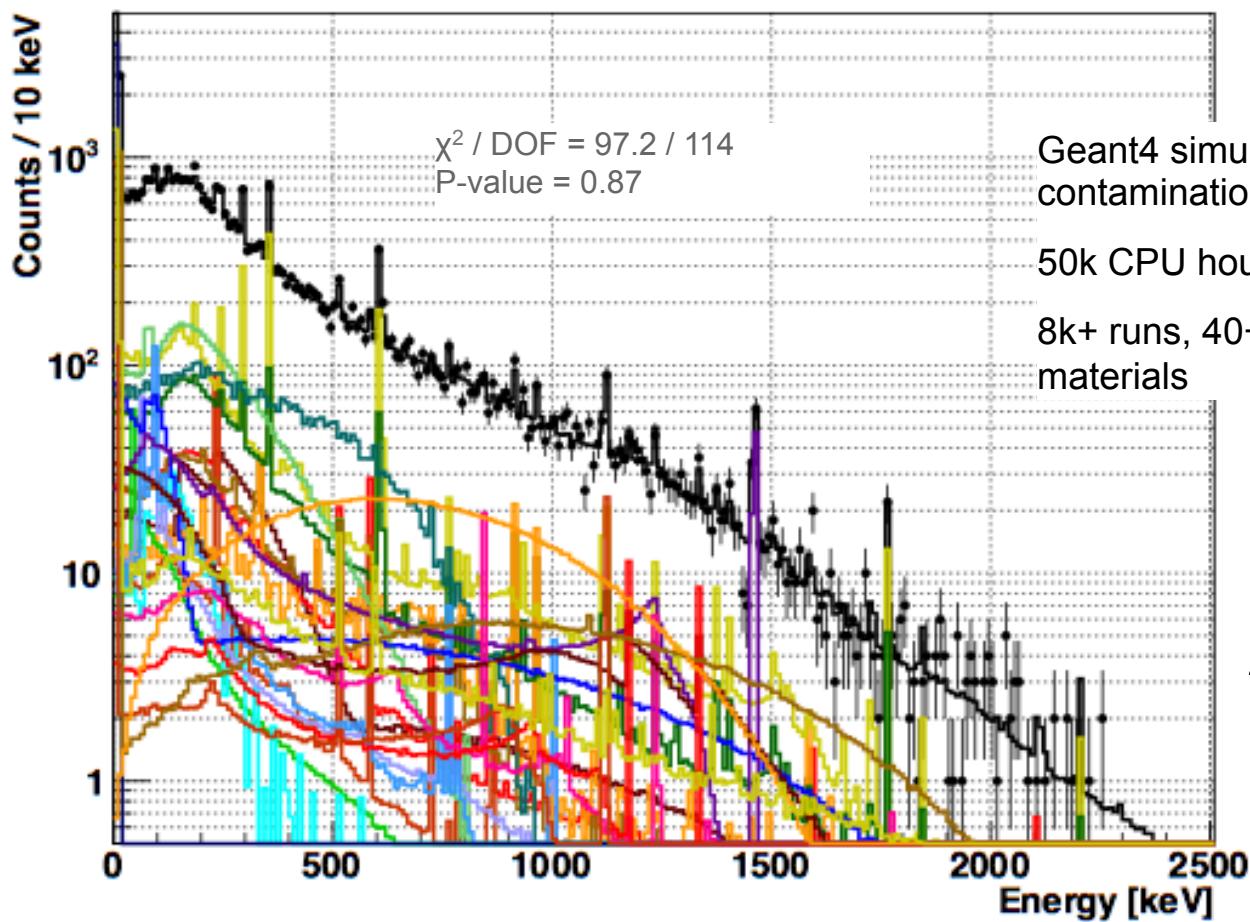
Natural radioactivity:

- **Pure material** (e.g. EFCu, clean plastic and others)
- **Shielding**
- **Analysis cuts** (PSD, granularity cuts)

Cosmogenic:

- **Deep underground** Combined efficiency of two layers of veto panel~99.9%,
Un-vetoed direct muon background<0.03 counts/ROI/t/y.
- **Muon veto**
- **Limit surface time of Ge** (shielded shipping and storage)
- **E-form Cu underground**
- **Analysis cuts** ^{68}Ge tag Single-Site Time Correlation Cut

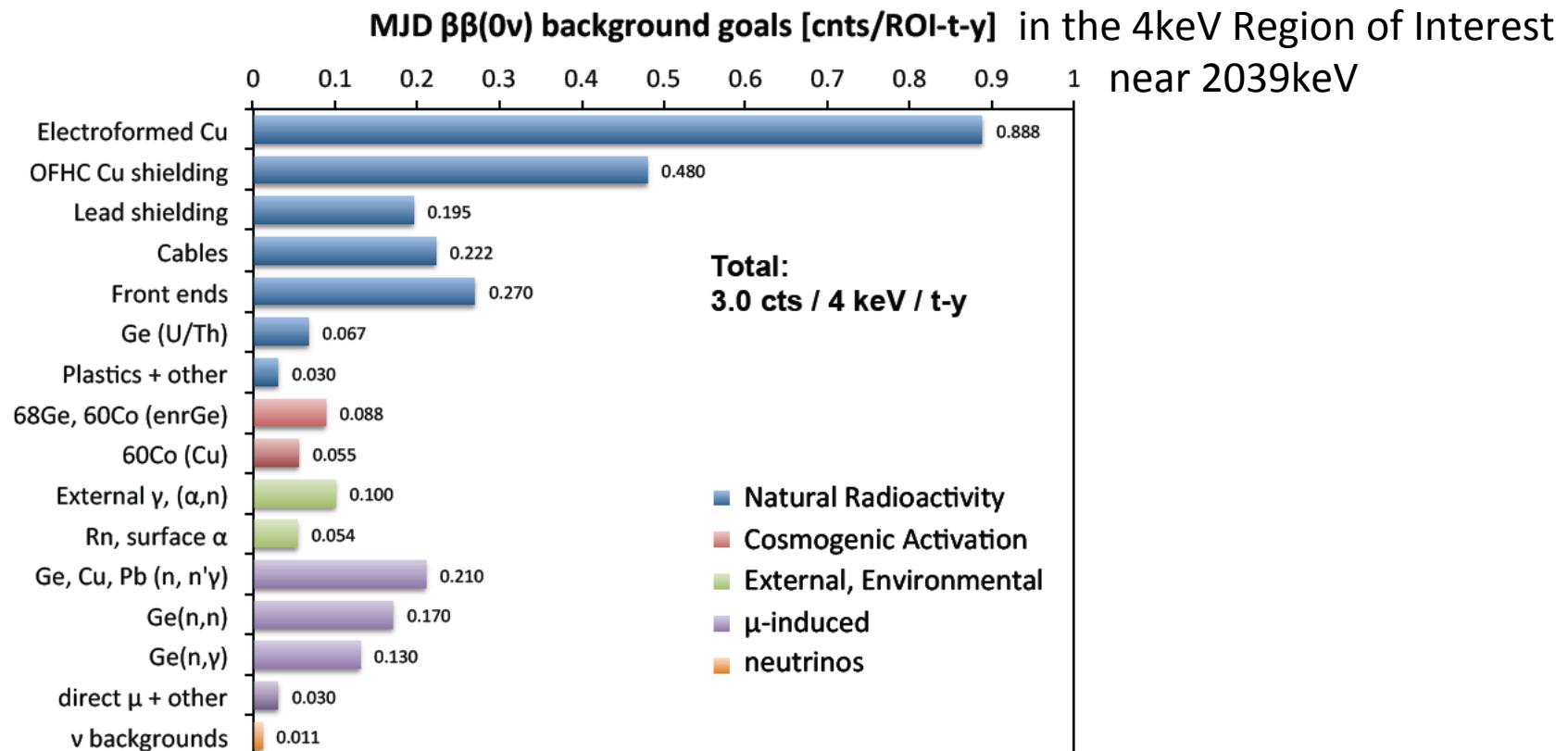
Background model fit of R&D Detector (MALBEK)





DEMONSTRATOR background budget

Based on assays of materials being used in MJD



2v $\beta\beta$ background is negligible due to excellent energy resolution

Simulated Background near $Q_{\beta\beta}$ after all cuts



Simulated spectra, 60 kg yrs, detector resolution + all cuts applied

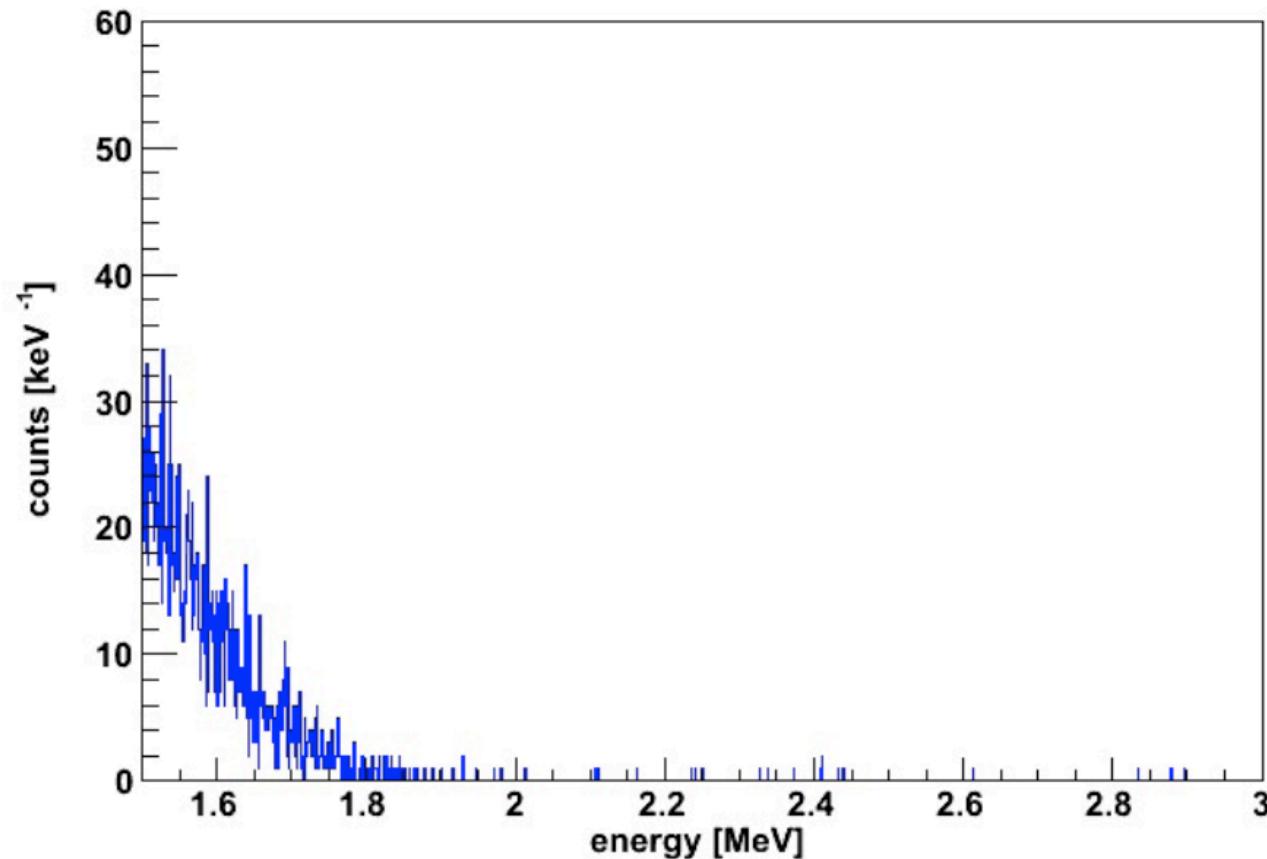


Figure adapted from
J.F.Wilkerson,
DOE ONP Comparative Review
June 25, 2013



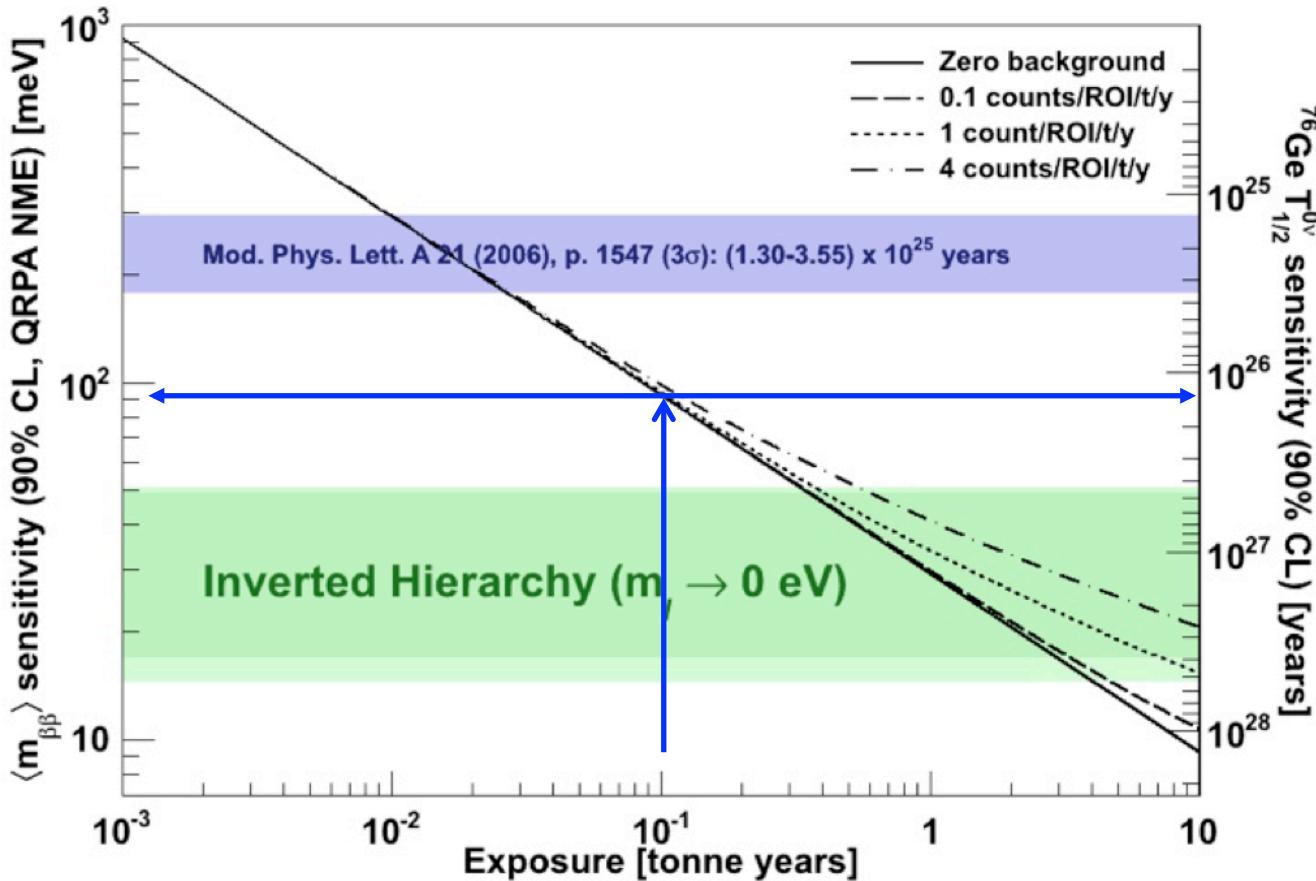
DEMONSTRATOR Status



- ✓ Infrastructure and cleanliness established
- ✓ Assayed all materials
- ✓ 75% required e-form copper produced
- ✓ 42.5 kg of 86% enriched 76Ge procured, refined to electronic grade with a 98% yield
- ✓ Accepted 10 enriched Ge det., 9.5kg in total
- ✓ Built two strings of natural Ge detectors built
- ✓ Fabricated prototype cryostat
- ✓ Built the associated vacuum system
- ✓ Shield construction in progress
- ✓ SlowControl and DAQ in use



DEMONSTRATOR schedule



Prototype Cryostat:
summer 2013

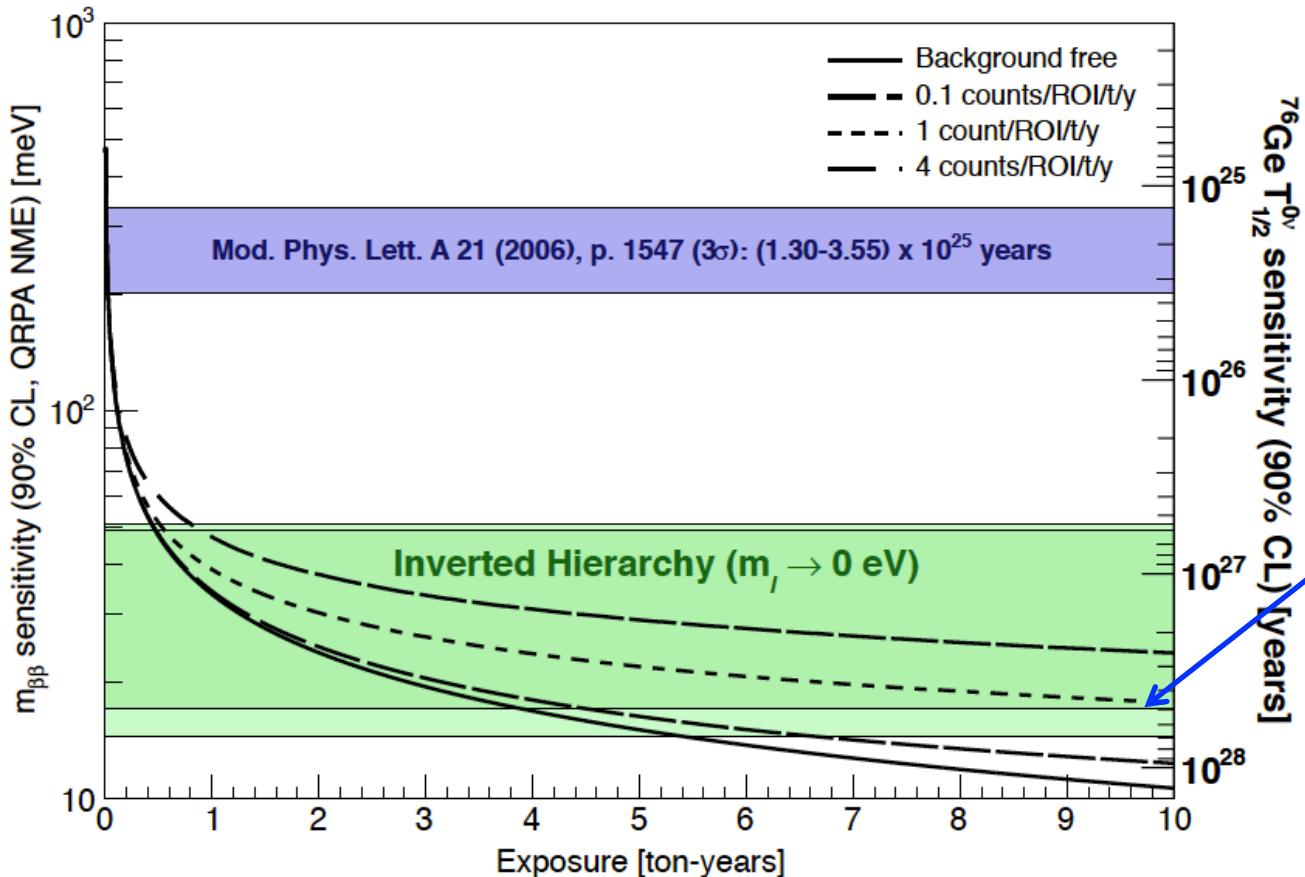
Cryostat 1:
end of 2013

Cryostat 2:
end of 2014

Run for 3 years,
exposure $\sim 100\text{kg}^*\text{y}$
Sensitive to $T \sim 10^{27}$ years



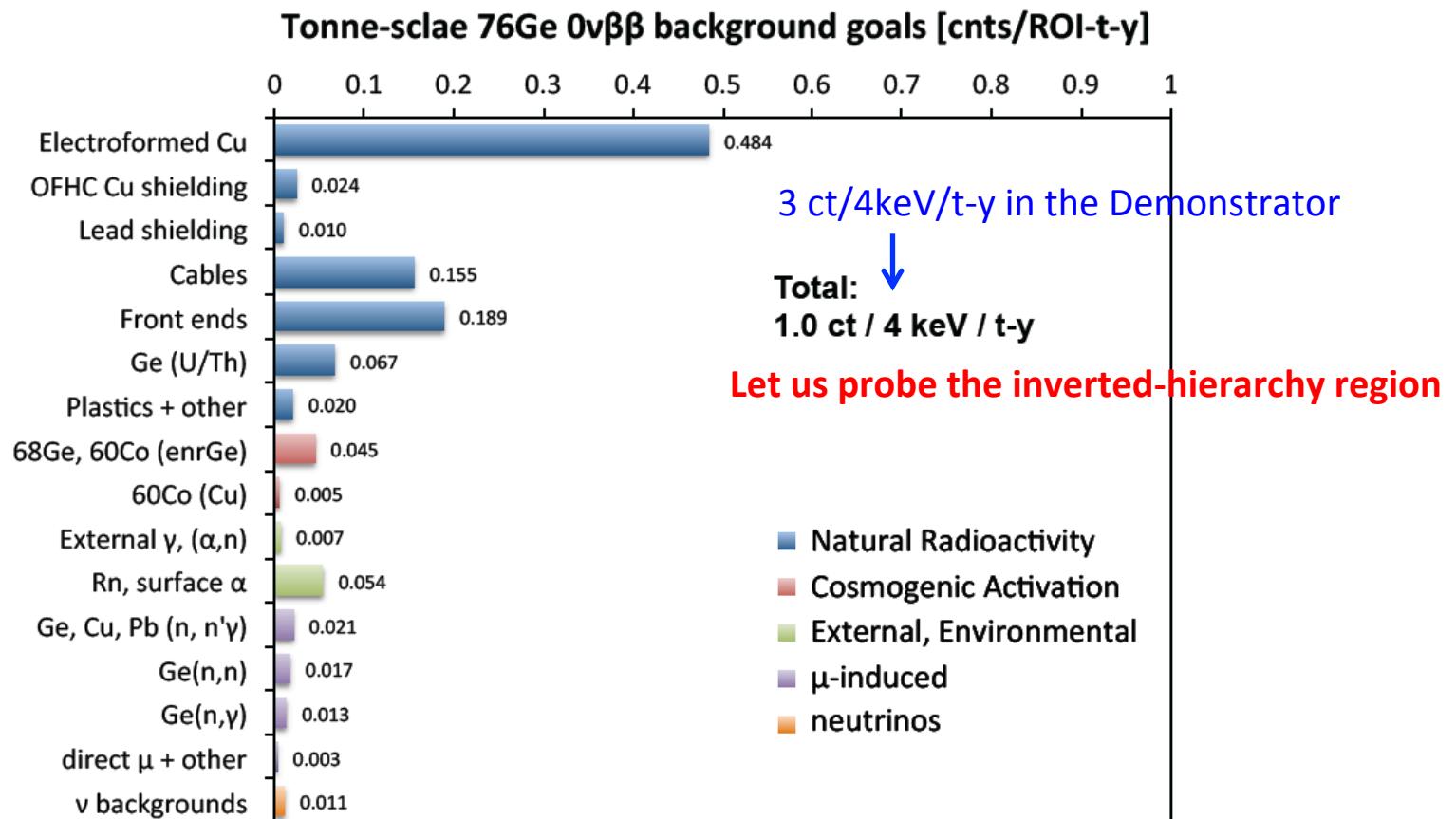
Tonne scale sensitivity



Background Projection for Tonne-scale



Scaling from MJD projects



The MAJORANA Collaboration



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Tom Burritt, Clara Cuesta, Jason Detwiler, Peter J. Doe, **Julieta Gruszko**, Greg Harper, **Jonathan Leon**, David Peterson, R. G. Hamish Robertson, Alexis Schubert, Tim Van Wechel



The END



THANK YOU!

backup

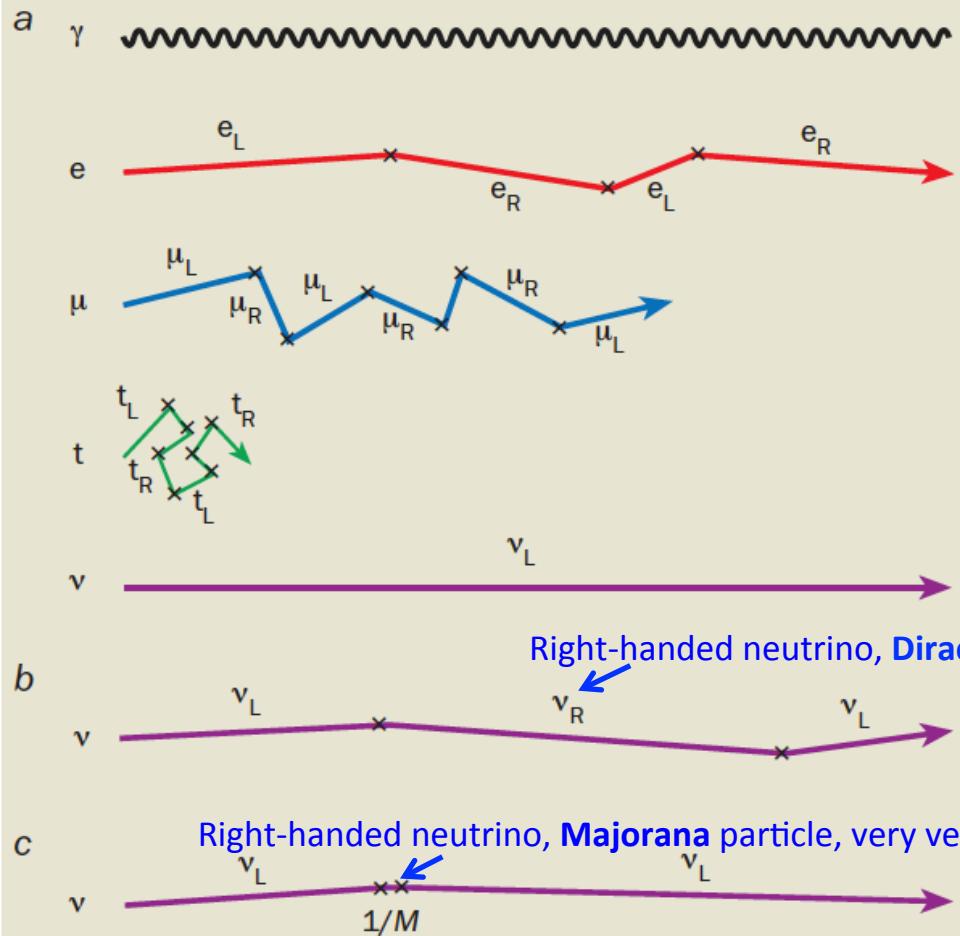


Seesaw mechanism: one motivation for majorana neutrinos



2 Neutrinos meet the Higgs boson

H. Murayama, Physics World, May 2002



Photons do not interact with Higgs, no mass

Other particles collider with Higgs, flipping the handedness, acquiring masses

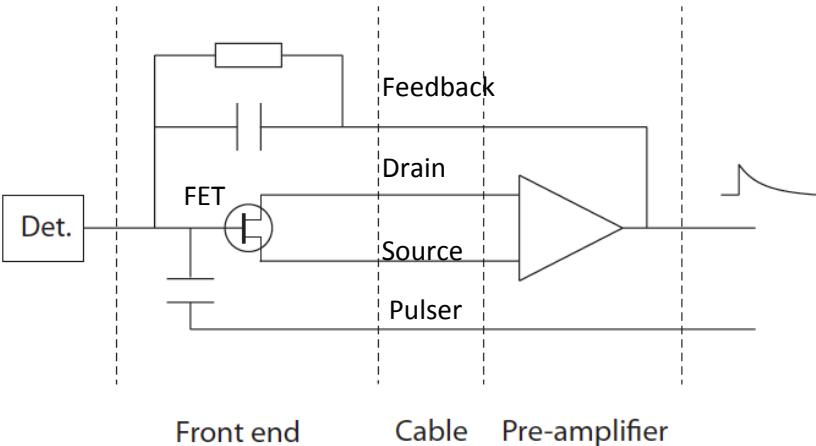
No right-handed neutrino, so no coupling to H.
Neutrinos having zero mass. Excluded

Right-handed neutrino, Dirac particle, similar mass, very very weak coupling to Higgs
Neutrinos having Dirac masses.
Why the interactions of ν_R are so weak?

Right-handed neutrino, Majorana particle, very very heavy mass, can only exist here due to uncertainty principle
Neutrinos having Majorana masses.
No need for very very weak interactions

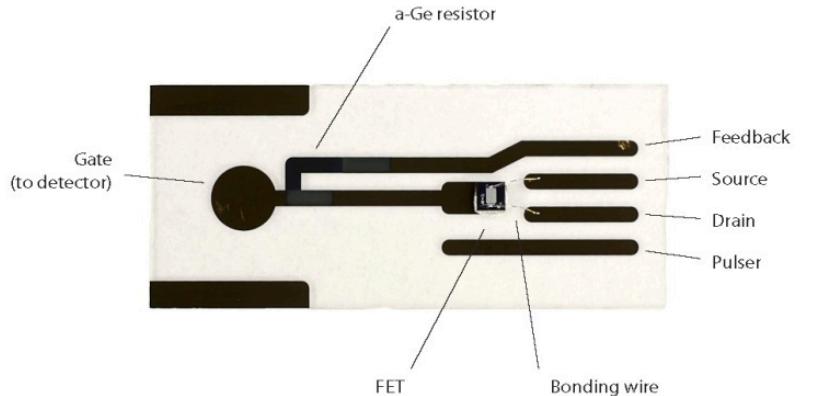


Low Mass Front End



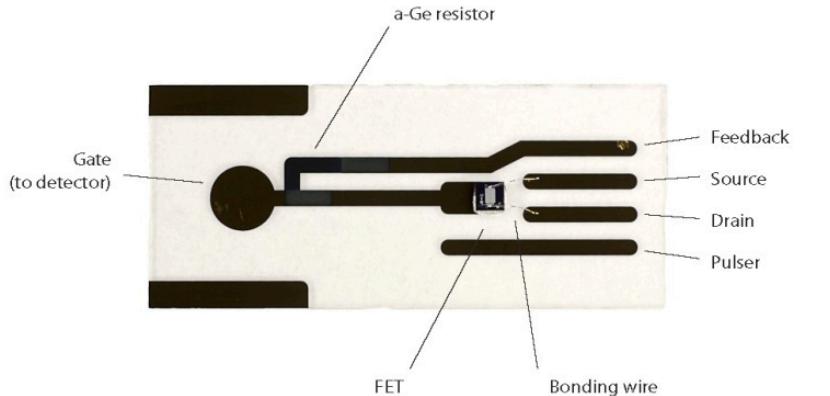
Self heating

Temperature can be Controlled by Drain to Source Voltage



- Fused silica substrate
- Au-Cr traces
- Amorphous-Ge resistor
- Low background
- Low noise

Low Mass Front End



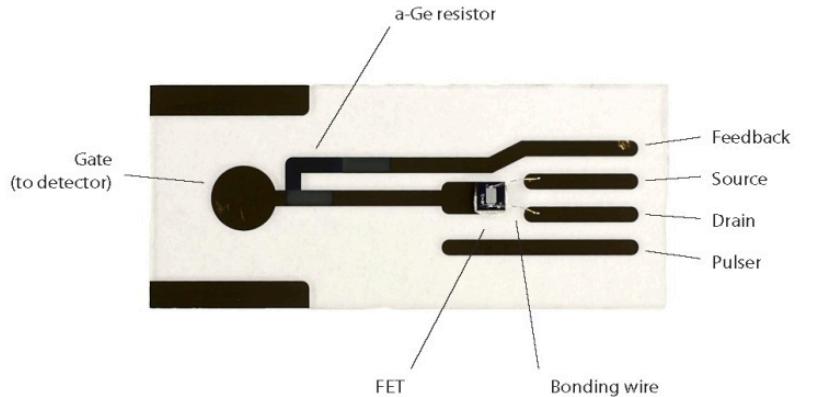
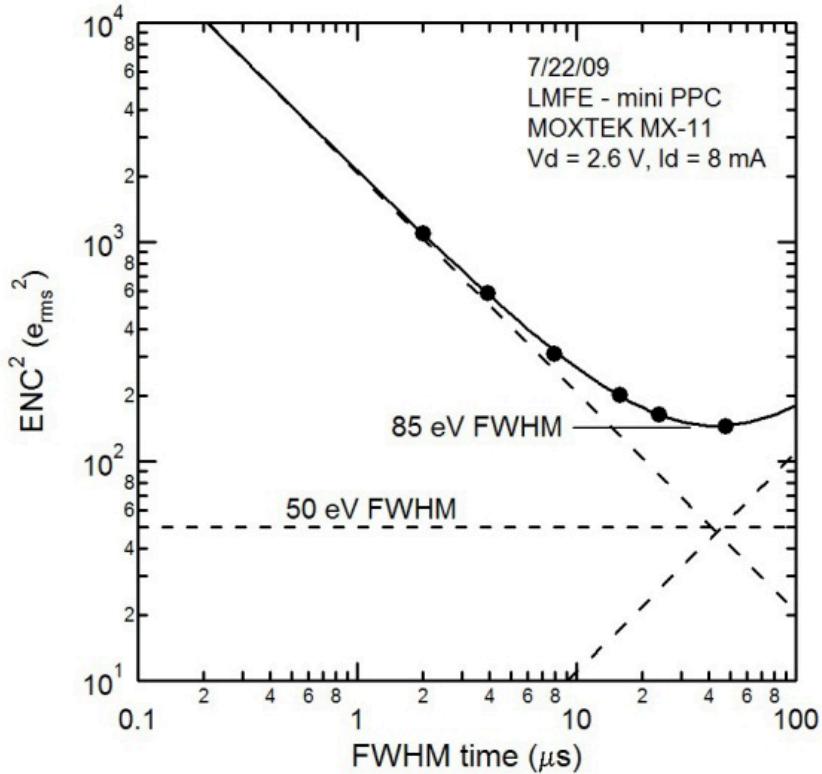
All material are selected to have low radioactivity

Component	Material	Purity (g/g)		Counts / ROI / t/y		Ref.
		^{232}Th	^{238}U	^{232}Th	^{238}U	
Substrate	Fused silica	101×10^{-12}	284×10^{-12}	0.0259	0.0616	MJ ICP-MS
Resistor	a-Ge	5×10^{-9}	5×10^{-9}	0.0001	0.0001	MJ ICP-MS
Traces	Au	$47(1) \times 10^{-9}$	$2.0(0.3) \times 10^{-9}$	0.0421	0.0015	MJ ICP-MS
Traces	Ti	$< 400 \times 10^{-12}$	$< 100 \times 10^{-12}$	~ 0	~ 0	MJ ICP-MS
FET	FET die	$< 2 \times 10^{-9}$	$< 141 \times 10^{-12}$	< 0.0107	< 0.0006	MJ ICP-MS
Bonding wire	Al	$91(2) \times 10^{-9}$	$9.0(0.4) \times 10^{-12}$	0.0004	~ 0	MJ ICP-MS
Epoxy	Silver epoxy	$< 70 \times 10^{-9}$	$< 10 \times 10^{-9}$	< 0.0685	< 0.0082	MJ gamma
Total				< 0.1476	< 0.0720	

- Fused silica substrate
- Au-Cr traces
- Amorphous-Ge resistor
- **Low background**
- Low noise



Low Mass Front End

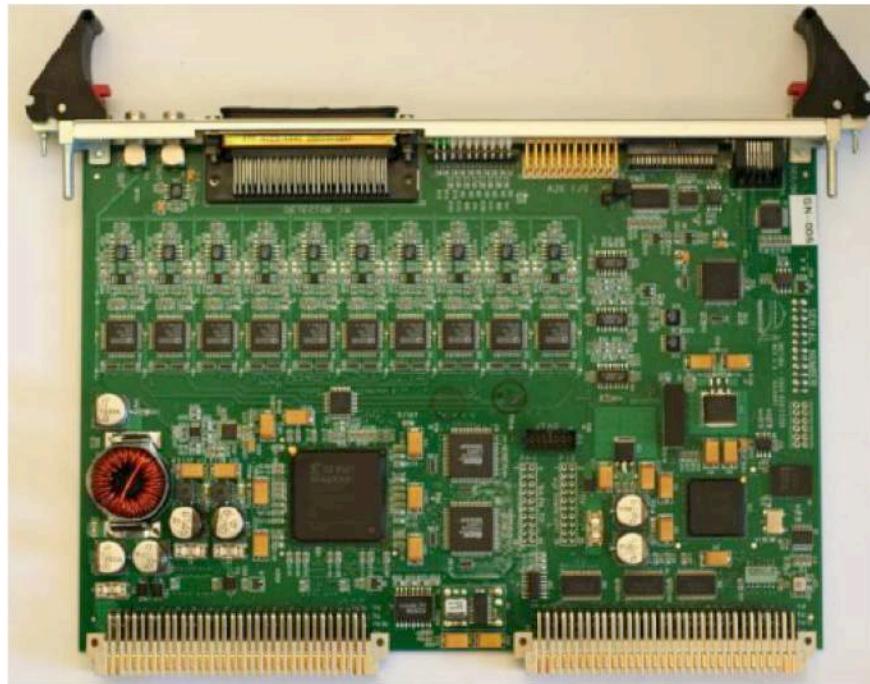


Equivalent noise:
55eV FWHM
without detector

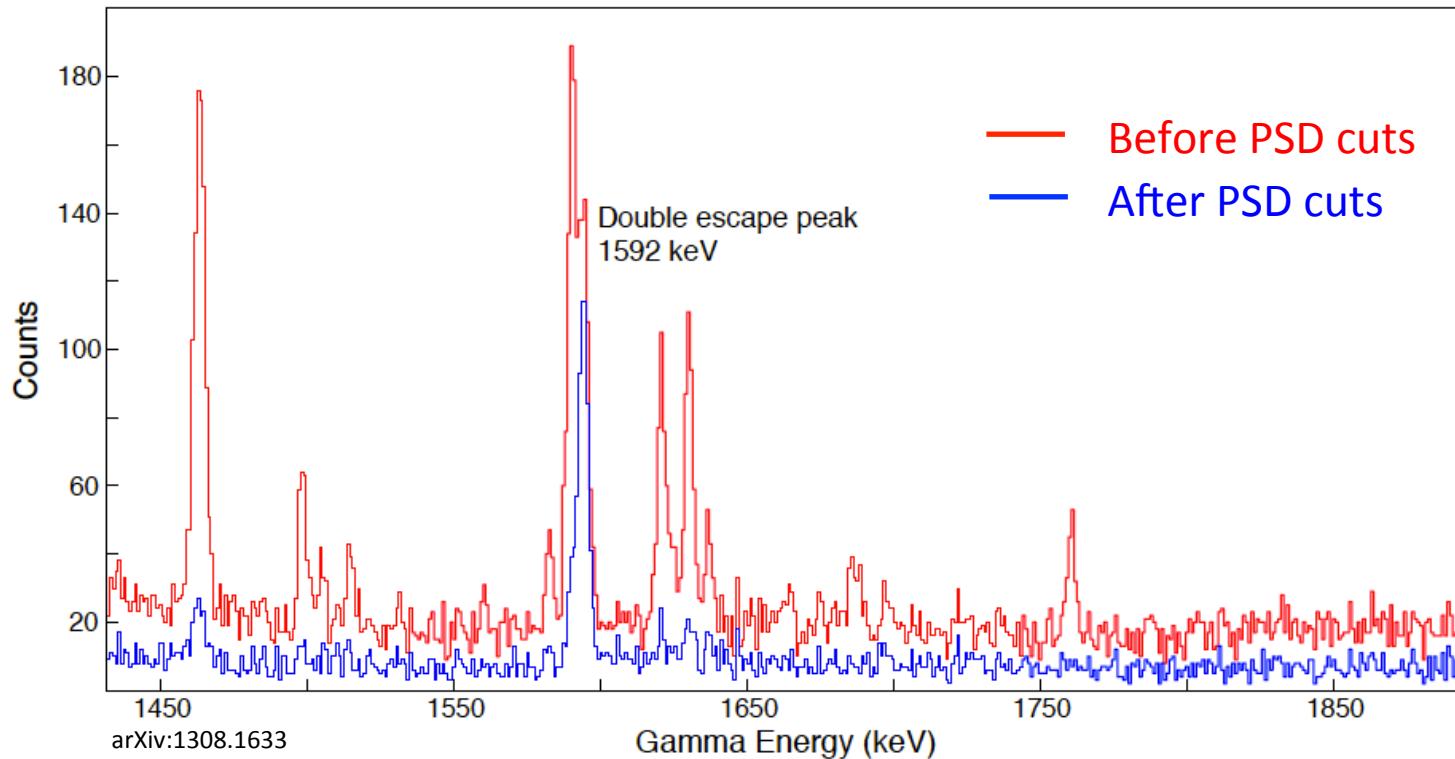
85eV FWHM
with a small detector



GRETINA digitizer card



Pulse-Shape-Discrimination





The most recent arguments

Why is the Conclusion of the GERDA Experiment Wrong ?

H.V. Klapdor-Kleingrothaus^a, I.V. Krivosheina^b
and S.N. Karpov^c

^a Heidelberg, Germany,
^b Heidelberg, Germany and Nishnij Novgorod, Russia,
^c JINR, Dubna, Russia

August 13, 2013

Abstract

The first results of the GERDA double beta experiment in Gran Sasso were recently presented. They are fully consistent with the Heidelberg-Moscow experiment, but *because of its low statistics cannot proof anything at this moment*. It is no surprise that the statistics is still far from being able to test the signal claimed by the Heidelberg-Moscow (HM) experiment. The energy resolution of the coaxial detectors is a factor of 1.5 worse than in the HM experiment. The *original goal* of background reduction to 10^{-2} counts/kg y keV, or by an order of magnitude compared to the Heidelberg-Moscow experiment, *has not been reached*. The background is *only* a factor 2.3 lower if we refer it to the experimental line width, i.e. in units counts/kg y energy resolution.

With pulse shape analysis (PSA) the background in the HM experiment around $Q_{\beta\beta}$ is 4×10^{-3} counts/kg y keV [1], which is *a factor of 4 (5 referring to the line width) lower* than that of GERDA with pulse shape analysis.

arXiv:1308.2541

“The **background** model is oversimplified and **not yet adequate**.

It is not shown that the lines of their background can be identified.

GERDA has to continue the measurement further 5 years, until they can responsibly present an understood background.”

... need much larger statistics/runtime

And

The 2006 claim was not excluded



Effective mass formula

the sum. The prediction is insensitive to θ_{13} and δm_{21}^2 because they are small. Setting $\theta_{13} = 0 = \delta m_{21}^2$, the following relation between M_{ee} and Σ is obtained for both hierarchies [297]:

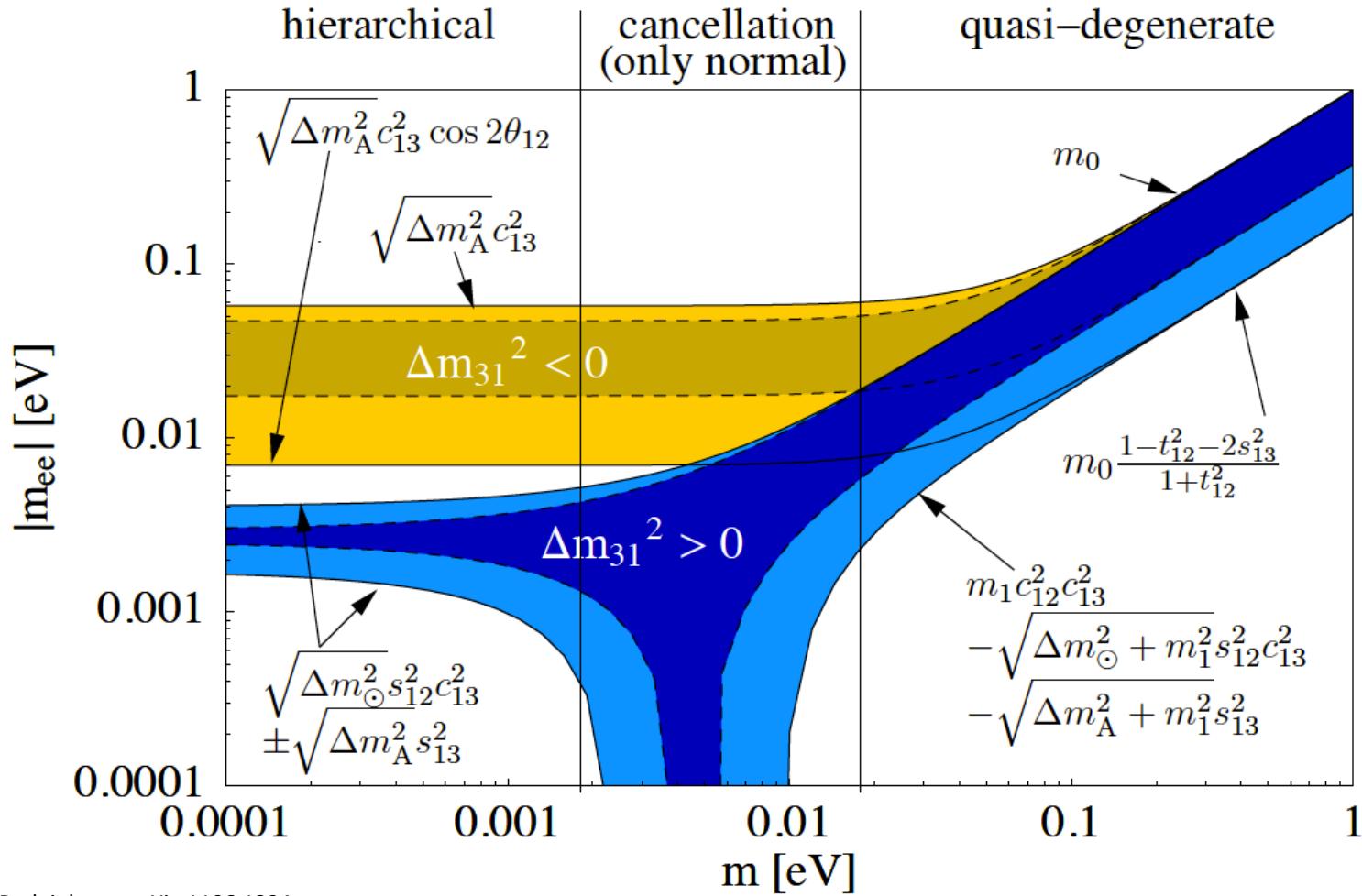
$$M_{ee} = \left(2\Sigma - \sqrt{\Sigma^2 + 3\delta m_{31}^2} \right) |c_{12}^2 + s_{12}^2 e^{i\phi}| / 3, \quad (7.8)$$

where ϕ is a Majorana phase. For a given measured value of M_{ee} both upper (since $\theta_{12} \neq \pi/4$) and lower bounds are implied for Σ . These bounds are displayed in figure 7.2. The present upper limit on M_{ee} is 0.35 eV at the 90% C.L. [301], with an overall factor of 3 uncertainty associated with the $0\nu\beta\beta$ nuclear matrix elements [302, 303]. A detection of neutrinoless double beta decay, corresponding to $M_{ee} = 0.39$ eV, has been reported [304], but this experimental result is highly controversial [305].

The physics of neutrinos, V Barger, D Marfatia, K Whisnant, 2012, Princeton University Press



The mass plot



Simulated Background near $Q_{\beta\beta}$



Simulated spectra, 60 kg yrs, detector resolution applied

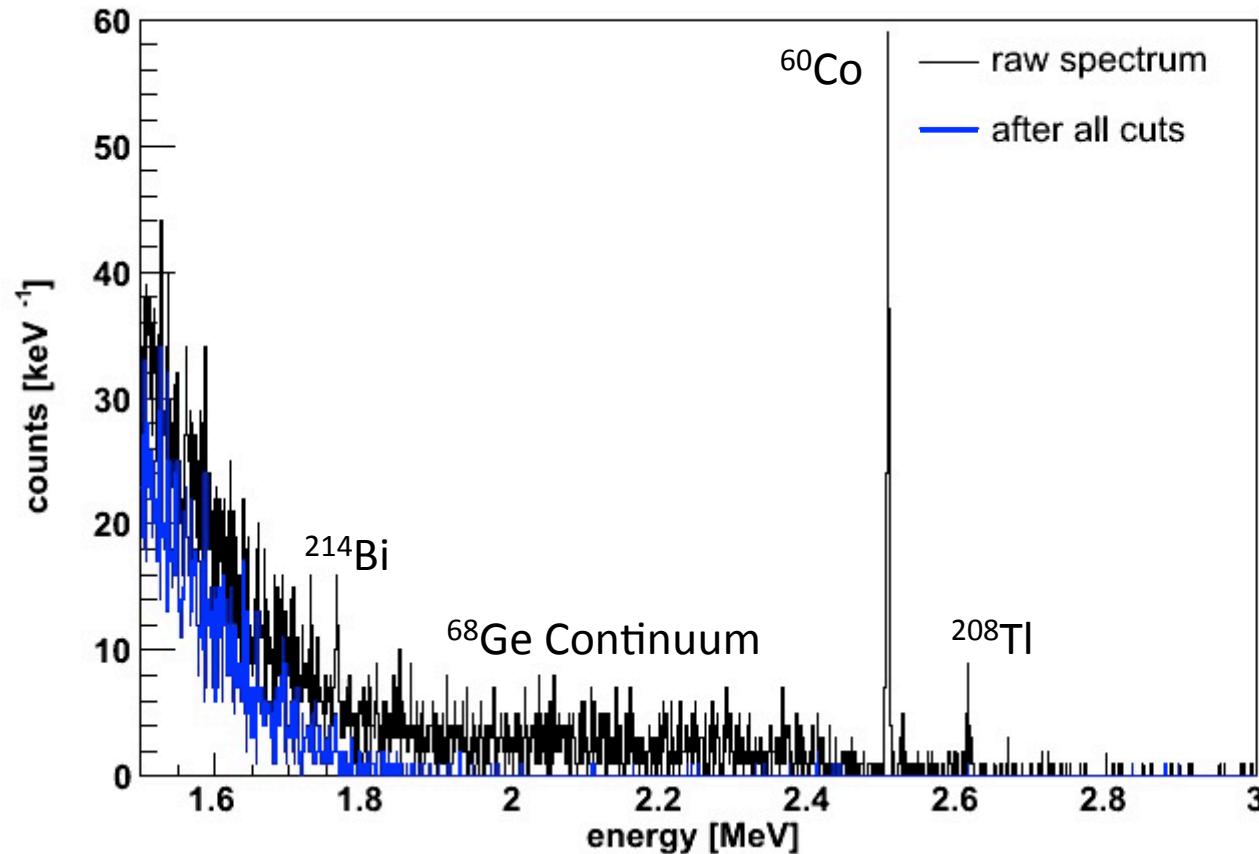


Figure adapted from
J.F.Wilkerson,
DOE ONP Comparative Review
June 25, 2013